

## Electronic Structure of Reduced Symmetry Peripheral Fused-Ring-Substituted Phthalocyanines

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Reduced symmetry phthalocyanines are finding use in an increasing number of industrial applications. A detailed understanding of the electronic structure of the  $\pi$ -system will greatly facilitate the design of new complexes, which fit the specifications required in many of these emerging high technology fields. NMR, electronic absorption, magnetic circular dichroism (MCD), and fluorescence emission and excitation spectra have been recorded for five generic metal phthalocyanine (MPc) derivatives in which additional benzene rings are fused either radially or obliquely onto at least one of the four peripheral benzo groups. The spectroscopy of four radially substituted compounds, zinc mononaphthotribenzotetraazaporphyrine (Zn3B1N), zinc monobenzotrinaphthotetraazaporphyrine (Zn1B3N), and two cis and trans zinc dibenzodinaaphthotetraazaporphyrine (Zn2B2N) isomers, is compared to that of the obliquely fused structural isomer of Zn3B1N (Zn3BoN) and the  $D_{4h}$  symmetry parent compounds, ZnPc and zinc naphthalocyanine (ZnNc). The selection of Zn(II) as the central metal eliminates the possibility of charge transfer between the metal and ring. None of the complexes studied contain any  $\sigma$ -bonded peripheral substituents.  $^1\text{H}$  NMR signals of the seven compounds are assigned on the basis of the coupling patterns, integrated proton numbers, and decoupling experiments. The SIMPFIT program was used to perform spectral band deconvolution analyses of absorption and MCD spectra. ZINDO molecular orbital calculations are described, and the optical spectra are assigned on the basis of the MO models that have been developed previously to account for the spectral properties of metal porphyrin (MP(-2)) and metal phthalocyanine (MPc(-2)) complexes.

### Introduction

Metal porphyrin complexes (MP(-2)) play a vital role in biological processes such as photosynthesis and respiration.<sup>1</sup> These complexes offer a unique chemistry that has several existing and potential industrial applications. The most commercially important group of the porphyrin class of molecules is the phthalocyanines, known systematically as tetrabenzotetraazaporphyrins.<sup>2</sup> Metal phthalocyanine complexes (MPc(-2)) have traditionally found use as dyes and pigments, and more recently as the photoconducting agent in photocopiers, because of their easy synthesis,

high stability, and the presence of intense  $\pi \rightarrow \pi^*$  transitions in the visible region.<sup>3</sup> In recent decades, there has been renewed interest in the use of phthalocyanine complexes in a variety of high technology fields including use in semi-

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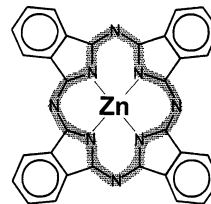
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conductor devices,<sup>4</sup> photovoltaic and solar cells,<sup>5</sup> electro-photography,<sup>6</sup> rectifying devices,<sup>7</sup> molecular electronics,<sup>8</sup> Langmuir–Blodgett films,<sup>9</sup> electrochromism in display devices,<sup>10</sup> low dimensional metals,<sup>11</sup> gas sensors,<sup>12</sup> liquid crystals,<sup>13</sup> and nonlinear optics,<sup>14</sup> and as photosensitizers<sup>15</sup> and electrocatalytic agents.<sup>16</sup> In a number of these applications, the wavelength of the major  $\pi \rightarrow \pi^*$  transitions in the UV–vis region is of critical importance. Because it is impossible to find an MPc(–2) complex that absorbs significantly for all desired wavelengths, there is considerable current interest in the electronic spectroscopy of structurally modified MPc(–2) molecules.<sup>17–22</sup> A detailed understanding



**Figure 1.** Molecular structure of a typical main group metal phthalocyanine showing the path of the 16-atom cyclic polyene ring that forms the basis of the 4-orbital calculations of Gouterman used to account for the two lowest energy major  $\pi \rightarrow \pi^*$  transitions of porphyrins and phthalocyanines.

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of the impact of different structural modifications on the electronic structure will greatly assist the development of new industrial applications for these molecules because it will facilitate the selection of previously uncharacterized molecules with a set of desired spectral properties via computer modeling.

In its simplest description, the optical spectroscopy of the  $\pi$ -system of MP(–2) and MPc(–2) can be accounted for by considering only the 16-atom, 18- $\pi$ -electron cyclic polyene aromatic ring that runs around the inner perimeter of the ligand, Figure 1. Gouterman's model, based on a 4-orbital linear combination of atomic orbitals (LCAO), has been widely used to describe the optical spectra of both MP(–2) and MPc(–2) complexes.<sup>3,23–26</sup> In an ideal 16-atom, 18- $\pi$ -electron cyclic polyene, the HOMO has an  $M_L$  value of  $\pm 4$  while the LUMO has an  $M_L$  value of  $\pm 5$ . This scheme predicts an allowed B transition ( $\Delta M_L = \pm 1$ ) and a forbidden Q transition ( $\Delta M_L = \pm 9$ ). MCD spectroscopy has confirmed the allowed and forbidden nature of the Q and B bands of MP(–2).<sup>27</sup> The  $A_1/D_0$  ratio (the MCD A term intensity divided by the absorption intensity) provides a direct spectral measurement of the angular momentum of the excited state. Gouterman's group demonstrated the validity of the 4-orbital approach through more detailed molecular orbital calculations.<sup>25,28</sup> Although Gouterman's models were developed over 30 years ago, no subsequent theoretical calculations have been reported that account for the optical spectra of MP(–2) and MPc(–2) complexes in a more satisfactory manner.<sup>29</sup> Gouterman's model has, therefore, been used as

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the theoretical framework within which to assign the bands observed in the optical spectra of a variety of MP(-2) and MPc(-2) complexes and their cation and anion radical species on the basis of spectral deconvolution studies of the UV-vis absorption and magnetic circular dichroism (MCD) spectroscopy<sup>3c,e,30</sup> using the SIMPFIT program.<sup>3a,e,31-36</sup> The major  $\pi \rightarrow \pi^*$  bands within the UV-vis region of MPc(-2) optical spectra are currently referred to as Q, B1, B2, N, L, and C in a sequence ordered by ascending energy on the basis of Gouterman's nomenclature.

The spectroscopy of structurally modified phthalocyanine complexes has received considerably less attention than that of MPc(-2) complexes with  $D_{4h}$  symmetry.<sup>3</sup> The impact of reducing the molecular symmetry through asymmetric  $\sigma$ -bonded substitution of the peripheral protons has been explored. The spectral changes observed are primarily related to the electron withdrawing or donating effects of the substituents on the  $\pi$ -system and are usually confined to slight shifts in the wavelengths of the major band centers.<sup>3,37</sup> In a recent paper, we studied the spectral changes that result from partial aza-nitrogen substitution at the 5, 10, 15, and 20 bridging positions of tetrabenzoporphyrin.<sup>38</sup> In this paper, we study the spectral impact of extending the  $\pi$ -system through asymmetric peripheral, fused benzene-ring substitutions. The extension of the  $\pi$ -system through peripheral aliphatic groups in naturally occurring porphyrin molecules such as protoporphyrin IX normally has only a minor impact on the optical spectroscopy. High molecular symmetry would still be inferred for these compounds on the basis of the MCD spectra because the peripheral substituents exert only a marginal influence on the spectroscopy of the inner perimeter cyclic polyene aromatic  $\pi$ -system.<sup>39</sup> Recent studies on the impact of ring expansion in the case of symmetric subzaporphyrin, subphthalocyanine, and subnaphthalocyanine com-

plexes have indicated that extending the  $\pi$ -system has a much more profound impact on the spectroscopy of the  $\pi$ -system.<sup>40</sup> Expansion of the heteroaromatic  $\pi$ -system by adding additional fused benzene rings offers the potential of maintaining the planarity of the  $\pi$ -system and a well-defined molecular symmetry while unambiguously modifying both the spectroscopic properties and electronic structure. Because many of the most industrially useful properties of MPc complexes are related directly to the planar heteroaromatic nature of the  $\pi$ -system, this structural modification should introduce new properties with considerable potential for use in new industrial applications.

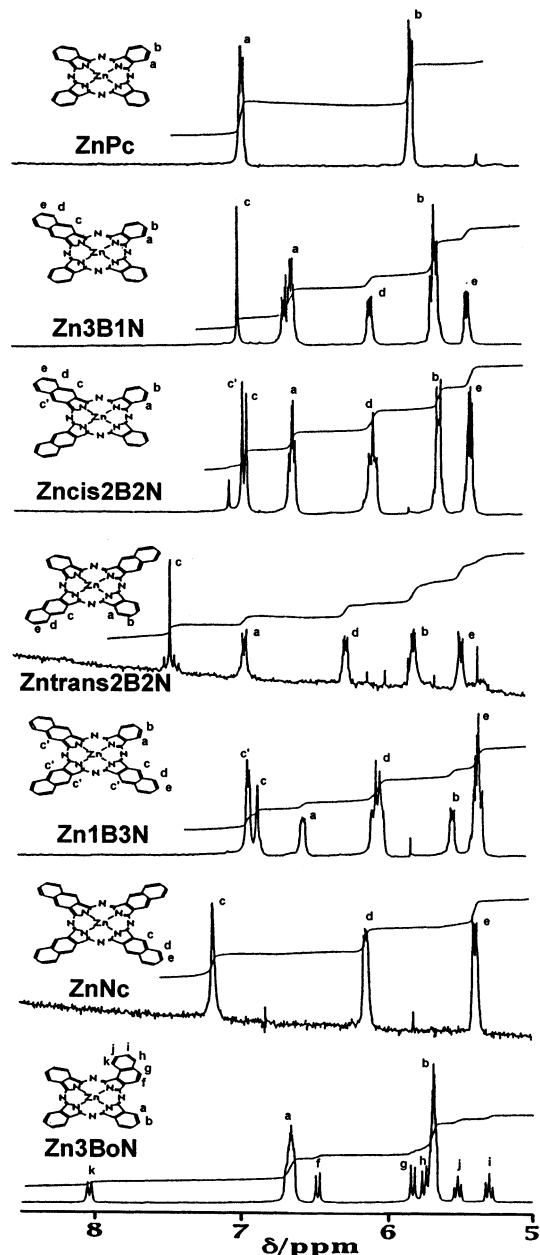
The molecular structures of the five asymmetric peripheral fused ring expanded complexes studied in this paper are shown in Figure 2. Zinc mononaphthotribenzotetraazaporphyrine (Zn3B1N) and monobenzotrinaphthotetraazaporphyrine (Zn1B3N) have  $C_{2v}(\text{III})$  symmetry, while the cis and trans isomers of dibenzodinaphthotetraazaporphyrine (Zn2B2N) have  $C_{2v}(\text{III})$  and  $D_{2h}$  symmetry, respectively. In  $C_{2v}(\text{III})$  symmetry, the  $C_2$  axis does not correspond to the main 4-fold axis of symmetry of MPc(-2) but lies instead within the plane of the  $\pi$ -system.<sup>41</sup> Zn3BoN is a lower symmetry structural isomer of Zn3B1N in which the additional outer fused ring is rotated obliquely out of alignment with the axes of symmetry of ZnPc resulting in  $C_s$  symmetry. NMR, electronic absorption, magnetic circular dichroism (MCD), and fluorescence emission and excitation spectroscopies are used to explore the impact of these structural modifications to the  $\pi$ -system. Because the complexes studied do not contain any peripheral substituent groups, <sup>1</sup>H NMR can be used to study the ring current generated by these five reduced symmetry heteroaromatic ligands. Zinc was selected as the central metal to eliminate ligand to metal and metal to ligand charge transfer bands. These spectra should, therefore, form a generic template for understanding the spectroscopy of this entire class of reduced symmetry phthalocyanine complexes.

## Experimental Section

(a) **Synthesis.** ZnPc(-2), Zn3B1N, Zncis2B2N, Zntrans2B2N, Zn1B3N, and ZnNc(-2) were synthesized by the mixed condensation method.<sup>18a</sup> Octabutoxylated complexes analogous to Zn3B1N, Zncis2B2N, Zntrans2B2N, and Zn1B3N have been reported previously.<sup>42</sup> Zn3BoN was obtained by mixing 1,2-dicyanonaphthalene (1 equiv), zinc acetate (1.1 equiv), hydroquinone (1 equiv), and phthalonitrile (3 equiv) and heating at 270 °C under nitrogen for 5 min. The residue was washed with water and methanol at room temperature and dissolved in THF for alumina column chromatography. The blue-green portion was isolated by gel-permeation chromatography using Bio-beads SX-8 (Bio-rad) and THF as eluent. Because the blue-green bands were not fully resolved, fractions were collected in a number of sample tubes and their electronic absorption spectra compared. By repeating this procedure several times, portions that gave a similar absorption spectrum were

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**Figure 2.**  $^1\text{H}$  NMR spectra and molecular structures of ZnPc(−2) **1**, Zn3B1N **2**, Zncis2B2N **3**, Zntrans2B2N **4**, Zn1B3N **5**, ZnNc(−2) **6**, and Zn3BoN **7**.

collected, concentrated, and recrystallized twice from THF. Analysis yielded C, 69.05; H, 3.11; N, 17.52, which is in good agreement with the anticipated C, 68.85; H, 2.89; N, 17.84. The parent  $m/z$  peak of 626 was observed by FAB mass spectrometry using a *m*-nitrobenzyl alcohol matrix.

**(b) Optical Spectroscopy.** UV–vis spectra were recorded with a Shimadzu UV-250 spectrophotometer, and magnetic circular dichroism (MCD) measurements were made with a JASCO J-720 spectrodichrometer. The JASCO electromagnet has a maximum field strength of 1.09 tesla (T). MCD intensity is expressed in terms of molar ellipticity per tesla,  $[\theta]_M = 10^4 \text{ deg mol}^{-1} \text{ dm}^3 \text{ cm}^{-1} \text{ T}^{-1}$ , and  $\Delta\epsilon_M$ , where  $[\theta]_M = 3300\Delta\epsilon_M$ . Fluorescence and excitation spectra were recorded with a Shimadzu RF-500 spectrofluorimeter with appropriate filters to eliminate scattered light. Fluorescence quantum yields ( $\Phi_F$ ) were determined by the comparative calibration method through the use of  $\text{H}_2\text{Pc}$  ( $\Phi_F = 0.60$ ) and ZnPc ( $\Phi_F = 0.30$ )<sup>43</sup> or 1,3,3,1',3',3'-hexamethylindotricarbocyanine ( $\Phi_F =$

0.28)<sup>44</sup> as standards. All solutions were purged with argon prior to fluorescence spectral measurement.

**(c) NMR Spectroscopy.**  $^1\text{H}$  NMR spectra were recorded with a JEOL GX-500 500 MHz spectrometer in dimethyl sulfoxide- $d_6$  with tetramethylsilane (TMS) as the standard. Because of the low solubility of many of the complexes studied (particularly Zntrans-2B2N and ZnNc(−2)) due to the absence of peripheral substituents, the signal had to be accumulated by averaging between 6000 and 48000 scans. The  $^1\text{H}$  signals obtained were assigned on the basis of the integrated intensities, coupling patterns, and the results of decoupling experiments, Figure 2.

**(d) MCD Spectroscopy.** MCD spectroscopy provides the ground and excited state degeneracy information, which is required to fully understand the electronic structure of high symmetry molecules. The additional information provided by the MCD technique is derived from three highly characteristic spectral features, the Faraday A, B, and C terms.<sup>30</sup> The derivative-shaped Faraday A term is temperature independent and identifies degenerate excited states, while the normally Gaussian-shaped C term is highly temperature dependent and identifies an orbitally degenerate ground state. Gaussian-shaped, temperature independent B terms arise from field-induced mixing between closely related states linked by a magnetic dipole transition moment. The ground states of divalent, main group  $D_{4h}$  MPc(−2) complexes are nondegenerate ( $^1A_{1g}$ ), and the accessible  $\pi \rightarrow \pi^*$  excited states are orbitally degenerate ( $^1E_u$ ). The MCD spectra are therefore dominated by intense  $x/y$ -polarized derivative-shaped A terms, Figure 3. Vibronically coupled and  $n \rightarrow \pi^*$  states can also transform as  $^1A_{2u}$ , giving rise to  $z$ -polarized B terms. Complexes with symmetry lower than  $D_{4h}$  give rise to oppositely signed coupled B terms usually with  $x$ - and  $y$ -polarization rather than A terms in the MCD spectra. The ordering of the sign of the A terms to high energy is not related to the polarization of the bands but to the orbital angular momentum associated with the ground and excited states in the absence of the perturbation to the  $D_{4h}$  symmetry of the inner cyclic polyene.<sup>24</sup> The ground state of Zntrans2B2N is  $^1A_g$ , and the accessible  $\pi \rightarrow \pi^*$  excited states are  $^1B_{3u}$  and  $^1B_{2u}$  under  $D_{2h}$  symmetry. In the case of Zn3B1N, Zncis2B2N, and Zn1B3N, the ground state is  $^1A_1$ , and the accessible  $\pi \rightarrow \pi^*$  excited states are  $^1A_1$  and  $^1B_1$  under  $C_{2v}(\text{III})$  symmetry giving rise to B terms of  $x$ - and  $z$ -polarization. The ground and excited states of the  $C_s$  symmetry Zn3BoN complex are  $^1A'$ .

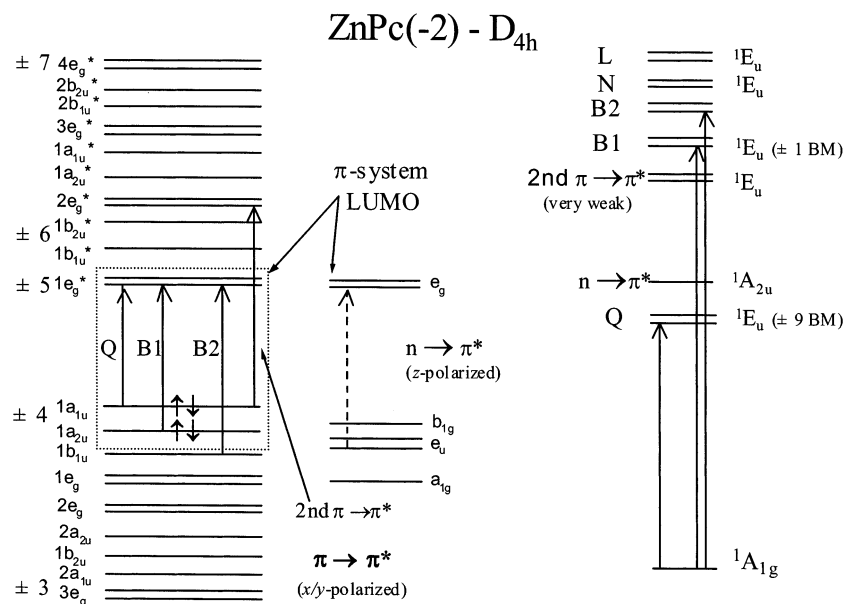
**(e) Band Fitting using SIMPFIT.** Band deconvolution has the potential to provide quantitative parameters for the allowed electronic transitions in extensively overlapped spectral envelopes. The major problem associated with band deconvolution analyses is that many different potential fits can usually be made for a single spectral data set. The unique approach developed during the design of the SIMPFIT program<sup>3e,45,46</sup> overcomes this problem through the simultaneous spectral deconvolution of the MCD and absorption spectral data sets recorded on the same solution. The MCD signal arises from the same electronic transitions as the UV–vis absorption spectrum, but the selection rules are different because the intensity mechanism depends on the magnetic dipole moment in addition to the electric dipole moment. The approach used in the SIMPFIT program rests on the fact that under the rigid shift assumption it can be assumed that the application of a magnetic field will make no difference to the spectral band shape function.<sup>30</sup> Fitting two

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**Figure 3.** Molecular orbital (A) and state level (B) diagrams of MPC(-2) showing the one-electron transitions that are predicted to give rise to absorption bands in the 280–1000 nm range. The orbital ordering on the left is based on Gouterman's model of the inner ring cyclic polyene.<sup>23</sup> The orbitals on the right are the four aza-nitrogen lone pair orbitals. The Q, second  $\pi \rightarrow \pi^*$ , B1, B2, N, and L bands have been identified through spectral deconvolution studies<sup>3a,e,31–36</sup> on the basis of Gouterman's SCMO-PPP-CI model.<sup>28</sup> The ordering of the MOs is based on more recent ZINDO calculations.<sup>54</sup> Transitions that give rise to  $x/y$ -polarized bands are represented with solid lines. Dashed lines are used for  $z$ -polarized transitions. The  $n\pi^*$  excited state is placed between the Q and B1  $\pi\pi^*$  excited states on the basis of a spectral deconvolution analysis of  $(\text{CN}^-)\text{ZnPc}(-2)$ .<sup>36</sup>

spectra recorded on the same solution with different band morphologies on the basis of two separate sets of selection rules with the same basic parameters removes much of the ambiguity that is normally associated with band deconvolution analyses. A new updated version of SIMPFIT<sup>45</sup> developed using Visual Basic was used during the band fitting reported in this paper.

**(f) ZINDO Molecular Orbital Calculations.** The structures for the ZINDO calculations, Figure 2, were developed through the use of a Fujitsu Inc. CAChe workstation.<sup>47</sup> The structures were then refined using a modified MM2 force field calculation in the Mechanics program of the CAChe system.<sup>48</sup> The structures were optimized at the restricted Hartree–Fock self-consistent field (SCF) level<sup>49</sup> using the ZINDO program in the CAChe software package.<sup>50,51</sup> These SCF optimizations were carried out at the intermediate neglect of differential overlap/1<sup>52</sup> (INDO/1) level of approximation. The structures were then used to calculate the UV–vis absorption spectra shown in Tables 1–3 and Tables S1–S3 (Supplementary Information). The active orbitals in the ZINDO/s configuration interaction (CI) calculation were the 25 lowest energy unoccupied and 25 highest energy occupied MOs. The heteroaromatic  $\pi$ -systems of the five complexes studied, Figure 2, have a

well-defined planar structure. The results obtained from ZINDO geometry optimization are therefore expected to be very similar to those obtained from other commonly used techniques such as density functional theory (DFT).<sup>53</sup> This was confirmed by running DFT geometry optimizations in which the B88-PW91 GGA function was used with DZVP basis sets.

## Results and Discussion

**ZINDO Calculations.** The calculated energies of the  $\pi$ -system MOs of a  $D_{4h}$  16 atom 18  $\pi$ -electron cyclic polyene  $\text{ZnC}_8\text{N}_8\text{H}_{16}$  model complex (ZnCP), zinc tetraaza-porphyrin (ZnTAP), ZnPc, Zn3B1N, Zn3BoN, Zncis2B2N, Zntrans2B2N, Zn1B3N, and ZnNc(-2) in the -11 to 3 eV range are shown schematically in Table 4 and Figure 4. The distinctive nodal and antinodal patterns observed for ZnPc(-2) MOs are observed consistently over the entire series of complexes. The orbitals have, therefore, been ordered relative to the corresponding ZnPc(-2) orbital so that the one-electron transitions responsible for the bands in the ZINDO calculated spectra of each complex can be directly compared to the existing model developed for the high symmetry parent complex.<sup>54</sup> The energies of  $\sigma$ -MOs associated primarily with the aza- and pyrrole-nitrogen lone pair orbitals in the -11 to 3 eV range are shown in Table 5. The correlation of the  $b_{1u}$ ,  $b_{2u}$ ,  $b_{1g}$ , and  $b_{2g}$  orbitals is different in the case of Zncis2B2N from that of Zn3B1N and Zn1B3N because the 2-fold axis of symmetry corresponds to the  $c_2'$  and not the  $c_2''$  axes of ZnPc(-2) and the symmetry of these orbitals is primarily dependent on the symmetry operations associated with these axes.

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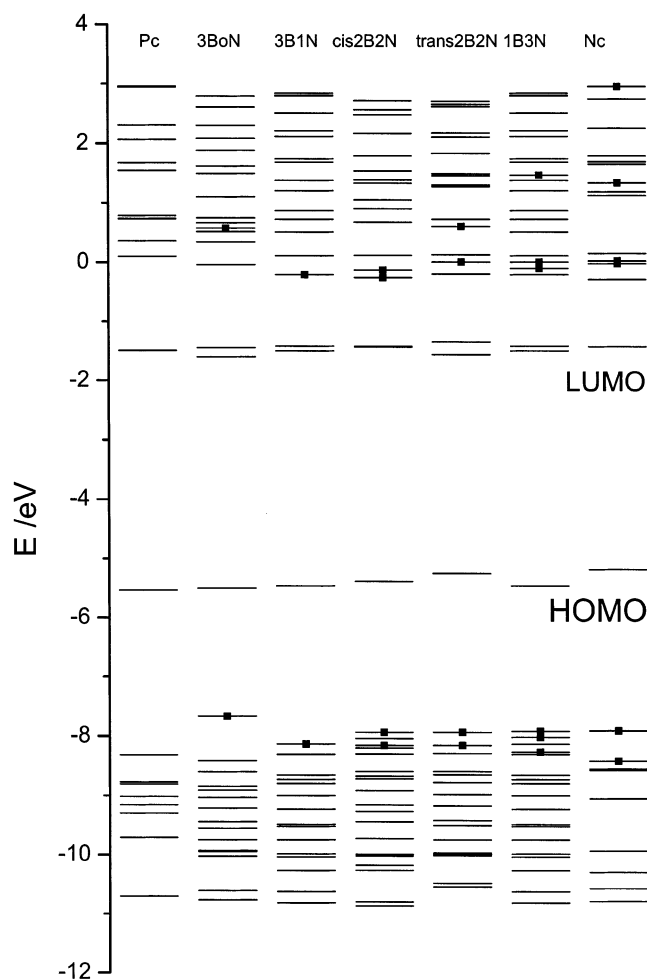
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**Table 1.** Calculated Electronic Excitation Spectrum of ZnPc(-2) Obtained from the ZINDO Program<sup>3c,e,54</sup>

no. <sup>a</sup>	sym <sup>b</sup>	calcd <sup>c</sup>	obsd <sup>d</sup>	wave function <sup>e</sup>	assignment <sup>f</sup>
1	<sup>1</sup> A <sub>1g</sub>				
2, 3	<sup>1</sup> E <sub>ux</sub>	14.8 (0.903)	14.8	-.950 1e <sub>g</sub> *←1a <sub>1u</sub> >+.260 1e <sub>g</sub> *←1a <sub>2u</sub> > +...	Q
6, 7	<sup>1</sup> E <sub>ux</sub>	30.0 (0.023)		-.978 2e <sub>g</sub> *←1a <sub>1u</sub> >+...	2nd π → π*
13	<sup>1</sup> A <sub>2u</sub>	34.1 (0.032)	16.5	.561 1e <sub>g</sub> *←e <sub>u</sub> <sup>N</sup> >-.472 1b <sub>2u</sub> *←b <sub>1g</sub> <sup>N</sup> > +...	n → π*
14, 15	<sup>1</sup> E <sub>ux</sub>	34.3 (0.428)	25.5	.808 1e <sub>g</sub> *←1b <sub>1u</sub> >-.255 1b <sub>2u</sub> *←1e <sub>g</sub> >- .221 1e <sub>g</sub> *←1a <sub>2u</sub> >+...	B1
16, 17	<sup>1</sup> E <sub>ux</sub>	34.7 (2.233)	29.9	-.815 1e <sub>g</sub> *←1a <sub>2u</sub> >-.348 3e <sub>g</sub> *←1a <sub>1u</sub> > -.252 1e <sub>g</sub> *←1b <sub>1u</sub> >+.228 1e <sub>g</sub> *←1a <sub>1u</sub> > +...	B2
21, 22	<sup>1</sup> E <sub>ux</sub>	36.5 (0.330)	33.6	-.574 3e <sub>g</sub> *←1a <sub>1u</sub> >+.371 1e <sub>g</sub> *←1a <sub>2u</sub> > -.259 1b <sub>1u</sub> *←1e <sub>g</sub> >-.251 2e <sub>g</sub> *←2a <sub>2u</sub> > -.222 2e <sub>g</sub> *←1b <sub>1u</sub> >+...	N
29, 31	<sup>1</sup> E <sub>ux</sub>	39.0 (0.014)	36.2	-.635 3e <sub>g</sub> *←1a <sub>1u</sub> >-.593 1e <sub>g</sub> *←1b <sub>2u</sub> > +.278 1e <sub>g</sub> *←2a <sub>1u</sub> >+.216 1b <sub>2u</sub> *←2e <sub>g</sub> > +...	L
34, 35	<sup>1</sup> E <sub>ux</sub>	41.1 (0.019)	40.7	-.744 1e <sub>g</sub> *←2a <sub>2u</sub> >+.451 1e <sub>g</sub> *←1b <sub>1u</sub> > -.293 1e <sub>g</sub> *←1b <sub>2u</sub> >+.255 3e <sub>g</sub> *←1a <sub>1u</sub> > +...	C

<sup>a</sup> The number of the state assigned in terms of ascending energy in the ZINDO calculation. Electronic transitions with a nonzero oscillator strength between 0 and 50000 cm<sup>-1</sup> are included in the table. <sup>b</sup> The symmetry of the state under *D*<sub>4h</sub> symmetry. <sup>c</sup> The calculated band energies (10<sup>3</sup>·cm<sup>-1</sup>) and oscillator strengths in parentheses. <sup>d</sup> Observed energies (10<sup>3</sup>·cm<sup>-1</sup>) from the spectral deconvolution data of Nyokong et al.<sup>31</sup> <sup>e</sup> The calculated wave functions based on the eigenvectors produced by the configuration interaction calculation of the ZINDO program. N denotes MOs associated with the aza-nitrogen lone pair orbitals. The orbital energies are shown in Table 4. <sup>f</sup> The assignment is described in the text.

**<sup>1</sup>H NMR Spectroscopy.** Analysis of the NMR of MPC(-2) complexes is based on the assumption that the <sup>1</sup>H signal shifts to lower field when there is a stronger ring current.<sup>37,55</sup> In tetra-*tert*-butylated H<sub>2</sub>TAP, H<sub>2</sub>Pc, H<sub>2</sub>Nc, and anthracocyanine, the pyrrole <sup>1</sup>H signals were observed at -2.47, -2.17, -1.64, and -0.73 ppm in CDCl<sub>3</sub>, respectively, because of the impact of ring expansion.<sup>19a,56</sup> Reduction of the molecular symmetry from *D*<sub>4h</sub> to either *C*<sub>2v</sub>(III) or *D*<sub>2h</sub> through the addition of peripheral fused benzene rings to MPC(-2) should therefore reduce the strength of the ring current of the inner cyclic polyene.<sup>57</sup> The anticipated shift of protons a and b to higher field can be seen in the NMR spectra of ZnPc(-2), Zn3B1N, Zncis2B2N, and Zn3BoN, and for protons c, d, and e in the spectra of Nc(-2) and Zn1B3N, Figure 2. When the π-system of the peripheral benzene rings in ZnPc(-2) is extended through the addition of naphthalene rings to form ZnNc(-2), the signal of proton b shifts to slightly higher field. The signals of protons a, c, and d do not appear to change appreciably as the number of naphthalene units is increased in the Zn3B1N, Zncis2B2N, and Zn1B3N series of *C*<sub>2v</sub>(III) complexes (the signals of protons c and c' split significantly in the case of Zncis2B2N and Zn1B3N). There is an abrupt shift to lower field in the signals of c, d, and e of Zntrans2B2N, which has *D*<sub>2h</sub> rather than *C*<sub>2v</sub>(III) symmetry, relative to those observed for the *D*<sub>4h</sub> ZnNc(-2) complex. In contrast, the signals of protons a and b are approximately the same as those in ZnPc(-2). Proton k of Zn3BoN sits at significantly lower field, because of the close proximity of one of the aza-nitrogens of the inner



**Figure 4.** Energy of the π-system MOs predicted by INDO/1 calculations for ZnPc(-2) (Pc), Zn3BoN, Zn3B1N, Zncis2B2N, Zntrans2B2N, and ZnNc(-2), Tables 4 and 5. Additional MOs associated primarily with the peripheral fused naphthalene units are indicated with a square.

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**Table 2.** Electronic Excitation Spectrum of Zn3BoN Obtained from the ZINDO Program<sup>d</sup>

no. <sup>b</sup>	sym <sup>c</sup>	calcd <sup>d</sup>	obsd <sup>e</sup>	wave function <sup>f</sup>	assignment <sup>g</sup>
1	A'				ground state
2	A'	14.5 (0.977)	14.7	<b>-.952 1a''*←1a''</b> )....	1a <sub>1u</sub> → 1e <sub>g</sub> * (Q)
3	A'	15.2 (0.803)	15.1	<b>-.945 2a''*←1a''</b> ) + ....	1a <sub>1u</sub> → 1e <sub>g</sub> * (Q)
6	A'	27.8 (0.133)		<u>-.754 2a''*←2a''</u> ) + .493 4a''*←1a'') + ....	RN → CP
8	A'	30.4 (0.316)		<u>.604 2a''*←2a''</u> ) + <b>.426 7a''*←1a''</b> ) -	RN → CP
				<b>.382 6a''*←1a''</b> ) + .367 5a''*←1a'') + ....	
9	A'	30.7 (0.380)		<u>-.725 6a''*←1a''</u> ) - <u>.550 2a''*←2a''</u> ) -	1a <sub>1u</sub> → 2e <sub>g</sub> * (2nd π → π*)
				<u>.250 9a''*←1a''</u> ) + ....	
11	A''	31.9 (0.005)		.678 1a''*←1a'' <sup>N</sup> ) + .447 2a''*←1a'' <sup>N</sup> ) -	n → π*
				.260 3a''*←1a'' <sup>N</sup> ) - .254 2a''*←3a'' <sup>N</sup> ) + ....	
12	A''	32.6 (0.001)		<u>-.674 2a''*←1a''<sup>N</sup></u> ) + .407 1a''*←1a'' <sup>N</sup> ) +	n → π*
				<u>.296 1a''*←3a''<sup>N</sup></u> ) - .211 2a''*←3a'' <sup>N</sup> ) + ....	
13	A'	32.9 (0.353)		<u>.612 8a''*←1a''</u> ) + <u>.354 1a''*←4a''</u> ) -	new "parity forbidden"
				<u>.276 2a''*←2a''</u> ) - <u>.214 2a''*←5a''</u> ) -	
				<u>.206 7a''*←2a''</u> ) + <b>.201 3a''*←4a''</b> ) + ....	
14	A'	33.4 (0.339)	25.6	<u>-.673 1a''*←6a''</u> ) + <u>.430 2a''*←5a''</u> ) +	1b <sub>1u</sub> → 1e <sub>g</sub> * (B2)
				<b>.358 1a''*←9a''</b> ) - <b>.267 1a''*←3a''</b> ) -	
				<u>.200 3a''*←9a''</u> ) + ....	
16	A'	34.0 (2.032)	27.1	<u>-.713 1a''*←3a''</u> ) - <u>.353 1a''*←4a''</u> ) +	1a <sub>2u</sub> → 1e <sub>g</sub> * (B1)
				<u>.283 1a''*←6a''</u> ) + <b>.242 10a''*←1a''</b> ) +	
				<b>.239 2a''*←1a''</b> ) - <u>.232 1a''*←2a''</u> ) + ....	
18	A'	34.6 (0.581)	28.1	<u>-.487 2a''*←3a''</u> ) - <b>.432 2a''*←6a''</b> ) +	1a <sub>2u</sub> → 1e <sub>g</sub> * (B1)
				<u>.363 1a''*←5a''</u> ) + <u>.267 9a''*←1a''</u> ) -	
				<u>.214 2a''*←4a''</u> ) + ....	
19	A'	34.7 (0.267)		<u>.538 8a''*←1a''</u> ) - <u>.504 1a''*←4a''</u> ) +	new "parity forbidden"
				<b>.303 1a''*←3a''</b> ) - <b>.261 2a''*←3a''</b> ) +	
				<b>.202 3a''*←4a''</b> ) + ....	
20	A'	36.0 (0.228)	29.6	<b>.578 2a''*←6a''</b> ) + <b>.339 2a''*←9a''</b> ) +	1b <sub>1u</sub> → 1e <sub>g</sub> * (B2)
				<u>.278 2a''*←7a''</u> ) + <u>.215 9a''*←1a''</u> ) -	
				<b>.208 3a''*←5a''</b> ) + ....	
21	A'	36.2 (0.538)	30.8	<b>.407 2a''*←3a''</b> ) + <u>.385 9a''*←2a''</u> ) -	B1/rN → CP
				<u>.327 1a''*←8a''</u> ) - <u>.257 4a''*←6a''</u> ) +	
				<b>.206 2a''*←9a''</b> ) + ....	
22	A'	36.3 (0.122)		<b>.506 10a''*←1a''</b> ) + <b>.272 1a''*←9a''</b> ) -	1a <sub>1u</sub> → 3e <sub>g</sub> * (B2)
				<b>.234 4a''*←5a''</b> ) - <u>.229 2a''*←8a''</u> ) -	
				<u>.227 1a''*←7a''</u> ) + <b>.222 1a''*←3a''</b> ) -	
				<b>.203 2a''*←10a''</b> ) + ....	
23	A'	36.4 (0.515)		<u>.399 2a''*←7a''</u> ) - <u>.393 1a''*←8a''</u> ) +	new "parity forbidden"
				<u>.360 1a''*←5a''</u> ) + <b>.360 2a''*←3a''</b> ) +	
				<u>.287 8a''*←1a''</u> ) - <b>.206 2a''*←9a''</b> ) + ....	
35	A'	40.3 (0.136)		<u>-.454 1a''*←9a''</u> ) + <b>.387 10a''*←1a''</b> ) -	2a <sub>2u</sub> → 1e <sub>g</sub> * (B2)
				<u>-.358 1a''*←7a''</u> ) + <u>.309 2a''*←4a''</u> ) -	
				<u>.302 8a''*←1a''</u> ) - <b>.238 1a''*←6a''</b> ) +	
				<u>.234 4a''*←10a''</u> ) + ....	
52	A'	45.7 (0.107)		<b>.856 14a''*←1a''</b> ) + <u>.387 17a''*←1a''</u> ) + ....	1a <sub>1u</sub> → 4e <sub>g</sub> * (B2)
53	A'	46.2 (0.131)		<u>-.723 16a''*←1a''</u> ) - <b>.387 15a''*←1a''</b> ) -	new "parity forbidden"
				<b>.288 1a''*←11a''</b> ) - <b>.275 2a''*←10a''</b> )	
				+ ....	
56	A'	47.1 (0.179)		<u>-.617 3a''*←3a''</u> ) - <b>.213 1a''*←11a''</b> )	new "parity forbidden"
				+ ....	
58	A'	47.4 (0.611)		<u>-.413 1a''*←11a''</u> ) + <u>.407 4a''*←2a''</u> ) -	1b <sub>2u</sub> → 1e <sub>g</sub> * (B2)
				<b>.267 14a''*←1a''</b> ) + <b>.254 3a''*←4a''</b> )	
				+ <u>.245 5a''*←4a''</u> ) + ....	
59	A'	48.0 (0.145)		<b>.428 2a''*←10a''</b> ) - <u>.398 3a''*←3a''</u> )	2a <sub>1u</sub> → 1e <sub>g</sub> * (B2)
				<u>-.212 15a''*←1a''</u> ) - <u>.211 16a''*←1a''</u> )	
				+ <u>.208 5a''*←3a''</u> ) + ....	
60	A'	48.7 (0.145)		<u>-.518 4a''*←2a''</u> ) + <u>.407 5a''*←2a''</u> )	RN → CP
				<u>-.268 7a''*←2a''</u> ) + <u>.239 8a''*←2a''</u> )	
				+ <u>.229 4a''*←3a''</u> ) + ....	
61	A'	49.2 (0.197)		<u>-.495 2a''*←11a''</u> ) + <u>.317 4a''*←3a''</u> )	1b <sub>2u</sub> → 1e <sub>g</sub> */new "parity forbidden"
				<u>-.307 3a''*←4a''</u> ) + <u>.257 5a''*←4a''</u> )	
				<u>-.221 5a''*←3a''</u> ) - <u>.214 3a''*←3a''</u> ) + ....	
62	A'	49.7 (0.442)		<u>-.459 2a''*←11a''</u> ) + <u>.271 1a''*←13a''</u> )	1b <sub>2u</sub> → 1e <sub>g</sub> * (B2)
				+ <b>.208 3a''*←4a''</b> ) - <u>.207 1a''*←12a''</u> )	
				<u>-.201 17a''*←1a''</u> ) + <u>.201 13a''*←2a''</u> )	
				+ ....	

<sup>a</sup> The energies of the molecular orbitals are shown in Tables 4 and 5, and spectral deconvolution parameters are reported in Tables S4 and S9. <sup>b</sup> The number of the state assigned in terms of ascending energy in the ZINDO calculation. All electronic transitions with an oscillator strength of greater than 0.1 between 0 and 50000 cm<sup>-1</sup> are included in the table. <sup>c</sup> The symmetry of the state assuming C<sub>s</sub> molecular symmetry. <sup>d</sup> The calculated band energies (10<sup>3</sup>·cm<sup>-1</sup>) and oscillator strengths in parentheses. <sup>e</sup> The band center energies (10<sup>3</sup>·cm<sup>-1</sup>) and intensities obtained from spectral deconvolution analysis are reported in Tables S4 and S9. <sup>f</sup> The calculated wave functions based on the eigenvectors produced by the configuration interaction calculation of the ZINDO program.<sup>47</sup> N denotes MOs associated with the aza-nitrogen lone pair orbitals. One-electron transitions shown in bold and italic are parity allowed and forbidden, respectively, in terms of the corresponding orbitals of ZnPc(-2) under D<sub>4h</sub> symmetry, Table 4. One-electron transitions that are underlined involve orbitals that are associated primarily with the new peripheral fused-ring extension to the π-system. <sup>g</sup> The assignment is based on the orbitals of ZnPc(-2) with the corresponding nodal and antinodal pattern, Tables 4 and 5.

**Table 3.** Electronic Excitation Spectrum of Zntrans2B2N Obtained from the ZINDO Program<sup>a</sup>

no. <sup>b</sup>	sym <sup>c</sup>	calcd <sup>d</sup>	obsd <sup>e</sup>	wave function <sup>f</sup>	assignment <sup>g</sup>
1	A <sub>g</sub>				ground state
2	B <sub>2u</sub>	13.8 (0.956)	13.7	- .957 1b <sub>2g</sub> *←1a <sub>u</sub> ⟩+....	1a <sub>1u</sub> → 1e <sub>g</sub> * (Q)
3	B <sub>3u</sub>	14.5 (1.184)	14.4	- .954 1b <sub>3g</sub> *←1a <sub>u</sub> ⟩+....	1a <sub>1u</sub> → 1e <sub>g</sub> * (Q)
10	B <sub>3u</sub>	31.8 (0.377)		- .573 1b <sub>3g</sub> *←2a <sub>u</sub> ⟩- .436 3b <sub>3g</sub> *←1a <sub>u</sub> ⟩+ <u>.282 1a<sub>u</sub>*←1b<sub>3g</sub>⟩</u> - .270 1b <sub>1u</sub> *←1b <sub>2g</sub> ⟩+ <u>.262 1b<sub>2g</sub>*←2b<sub>1u</sub>⟩+ .220 1b<sub>2g</sub>*←1b<sub>1u</sub>⟩+....</u> - .709 1b <sub>2g</sub> *←1b <sub>1u</sub> ⟩- .379 1b <sub>2g</sub> *←2b <sub>1u</sub> ⟩- <u>.373 1b<sub>2g</sub>*←3b<sub>1u</sub>⟩+ .0274 1b<sub>1u</sub>*←1b<sub>2g</sub>⟩</u> +....	2N → CP
12	B <sub>3u</sub>	33.4 (2.774)	26.2		1a <sub>2u</sub> → 1e <sub>g</sub> * (B1)
17	B <sub>3u</sub>	34.4 (0.707)	27.6	- .603 1b <sub>2g</sub> *←2b <sub>1u</sub> ⟩- .472 1b <sub>2g</sub> *←3b <sub>1u</sub> ⟩- <u>.461 1b<sub>2g</sub>*←1b<sub>1u</sub>⟩+ .227 1a<sub>u</sub>*←2b<sub>3g</sub>⟩</u> +....	1b <sub>1u</sub> → 1e <sub>g</sub> * (B2)
20	B <sub>2u</sub>	34.8 (0.767)	28.8	- .718 3b <sub>2g</sub> *←1a <sub>u</sub> ⟩- .328 1b <sub>3g</sub> *←1b <sub>1u</sub> ⟩+ <u>.319 4b<sub>2g</sub>*←1a<sub>u</sub>⟩- .281 1b<sub>3g</sub>*←3b<sub>1u</sub>⟩- .236 1b<sub>2g</sub>*←2a<sub>u</sub>⟩+....</u> <u>.650 1b<sub>2g</sub>*←1b<sub>2u</sub><sup>N</sup>⟩+ .468 1b<sub>3g</sub>*←1b<sub>3u</sub><sup>N</sup>⟩+ .383 1a<sub>u</sub>*←1b<sub>1g</sub><sup>N</sup>⟩+....</u>	1a <sub>1u</sub> → 2e <sub>g</sub> * (2nd π → π*)
22	B <sub>1u</sub>	35.3 (0.028)			n → π*
23	B <sub>2u</sub>	35.6 (0.924)	30.0	- .726 1b <sub>3g</sub> *←2b <sub>1u</sub> ⟩+ <u>.348 1b<sub>2g</sub>*←2a<sub>u</sub>⟩</u> - <u>.311 1b<sub>3g</sub>*←1b<sub>1u</sub>⟩+....</u>	1b <sub>1u</sub> → 1e <sub>g</sub> * (B2)
25	B <sub>2u</sub>	36.0 (0.648)	31.5	- .724 1b <sub>3g</sub> *←1b <sub>1u</sub> ⟩+ <u>.337 1b<sub>2g</sub>*←2a<sub>u</sub>⟩</u> - <u>.278 1b<sub>3g</sub>*←2b<sub>1u</sub>⟩+ .224 1b<sub>2g</sub>*←1a<sub>u</sub>⟩+....</u>	1a <sub>2u</sub> → 1e <sub>g</sub> * (B1)
32	B <sub>1u</sub>	39.0 (0.004)		- .695 1b <sub>3g</sub> *←1b <sub>3u</sub> <sup>N</sup> ⟩+ <u>.616 1b<sub>2g</sub>*←1b<sub>2u</sub><sup>N</sup>⟩</u> +....	n → π*
33	B <sub>2u</sub>	39.2 (0.475)	32.8	- .663 4b <sub>2g</sub> *←1a <sub>u</sub> ⟩- .458 1b <sub>2g</sub> *←3a <sub>u</sub> ⟩- <u>.312 1b<sub>1u</sub>*←1b<sub>3g</sub>⟩- .295 2b<sub>2g</sub>*←2a<sub>u</sub>⟩</u> <u>- .252 1b<sub>3g</sub>*←1b<sub>1u</sub>⟩+....</u>	1a <sub>1u</sub> → 3e <sub>g</sub> *
42	B <sub>2u</sub>	41.2 (0.723)		- .553 1b <sub>1u</sub> *←1b <sub>3g</sub> ⟩+ <u>.471 2b<sub>2g</sub>*←2a<sub>u</sub>⟩</u> - <u>.384 4b<sub>2g</sub>*←1a<sub>u</sub>⟩- .258 2b<sub>2g</sub>*←2a<sub>u</sub>⟩+....</u>	2N → CP
56	B <sub>3u</sub>	46.9 (1.134)		- .408 1a <sub>u</sub> *←1b <sub>3g</sub> ⟩- .396 1b <sub>1u</sub> *←1b <sub>2g</sub> ⟩+ <u>.294 3b<sub>1u</sub>*←1b<sub>2g</sub>⟩- .242 2b<sub>2g</sub>*←2b<sub>1u</sub>⟩</u> - <u>.238 3b<sub>3g</sub>*←2a<sub>u</sub>⟩+ .227 2b<sub>2g</sub>*←1b<sub>1u</sub>⟩+....</u>	2N → CP
59	B <sub>2u</sub>	47.4 (0.280)		<u>.581 1b<sub>2g</sub>*←3a<sub>u</sub>⟩- .367 2b<sub>1u</sub>*←2b<sub>3g</sub>⟩</u> + <u>.280 2b<sub>3g</sub>*←2b<sub>1u</sub>⟩+ .280 4b<sub>2g</sub>*←1a<sub>u</sub>⟩+ .252 2b<sub>3g</sub>*←1b<sub>1u</sub>⟩- .217 1b<sub>3g</sub>*←1b<sub>1u</sub>⟩+....</u>	1b <sub>2u</sub> → 1e <sub>g</sub> *
63	B <sub>3u</sub>	48.8 (0.264)		<u>.603 2b<sub>2g</sub>*←1b<sub>1u</sub>⟩+ .367 2b<sub>2g</sub>*←2b<sub>1u</sub>⟩</u> + <u>.293 1a<sub>u</sub>*←1b<sub>3g</sub>⟩+ .259 1b<sub>1u</sub>*←2b<sub>2g</sub>⟩</u> - <u>.238 2b<sub>2g</sub>*←3b<sub>1u</sub>⟩+ .223 1b<sub>1u</sub>*←1b<sub>2g</sub>⟩+....</u>	2N → CP

<sup>a</sup> The energies of the molecular orbitals are shown in Tables 4 and 5, and spectral deconvolution data are reported in Tables S7 and S12. <sup>b</sup> The number of the state assigned in terms of ascending energy in the ZINDO calculation. All electronic transitions with an oscillator strength of greater than 0.1 between 0 and 50000 cm<sup>-1</sup> are included in the table. <sup>c</sup> The symmetry of the state assuming D<sub>2h</sub> molecular symmetry. <sup>d</sup> The calculated band energies (10<sup>3</sup>·cm<sup>-1</sup>) and oscillator strengths in parentheses. <sup>e</sup> The band center energies (10<sup>3</sup>·cm<sup>-1</sup>) and intensities obtained from spectral deconvolution analysis are reported in Tables S7 and S12. <sup>f</sup> The calculated wave functions based on the eigenvectors produced by the configuration interaction calculation of the ZINDO program.<sup>47</sup> N denotes MOs associated with the aza-nitrogen lone pair orbitals. One-electron transitions that are underlined involve orbitals that are associated primarily with the new peripheral fused-ring extension to the π-system. <sup>g</sup> The assignment is based on the orbitals of ZnPc(-2) with the corresponding nodal and antinodal pattern, Tables 4 and 5.

cyclic polyene.<sup>58</sup> To interpret the NMR of reduced symmetry phthalocyanine derivatives, the ring currents associated with the 16-atom inner cyclic polyene, benzene, and naphthalene rings systems clearly all have to be taken into consideration.

**Optical Spectroscopy. (a) UV–Vis Absorption, Emission, and MCD Spectroscopy.** The UV–vis absorption and MCD spectra of ZnPc, Zn3B1N, Zn3BoN, Zncis2B2N, Zntrans2B2N, Zn1B3N, and ZnNc(-2) are shown in Figures 5 and 6, respectively. The A terms associated with the lowest energy π → π\* band of ZnPc and ZnNc(-2) are replaced by two intense, oppositely signed, close-lying B terms in the case of Zn3B1N, Zn3BoN, Zncis2B2N, Zntrans2B2N, and Zn1B3N. Michl<sup>24</sup> has demonstrated that the -ve/+ve ordering of the signs of a coupled B term to high energy in reduced symmetry cyclic polyene π-systems is consistent with an A term involving an excited state that has greater orbital angular momentum (OAM) than the ground state if D<sub>4h</sub> or higher symmetry had been maintained.<sup>24</sup>

The impact of perturbations on the electronic structure of the ideal cyclic polyene running around the inner perimeter of the porphyrin π-system can be predicted qualitatively by considering the location of the nodes and antinodes of the four frontier orbitals by the “Symmetry-adapted Perturbation Method”.<sup>18c,19b,23,24,59</sup> The pyrrolic nitrogens lie on a node of the HOMO a<sub>2u</sub> orbital and are very close to the nodes of the second HOMO a<sub>1u</sub> orbital of MP(-2). As a result, the addition of pyrrolic nitrogens to the cyclic polyene has almost no impact on the separation between the 1a<sub>1u</sub> and 1a<sub>2u</sub> HOMO orbitals (ΔHOMO) so these two orbitals can be viewed as being essentially degenerate in the case of MP(-2). The addition of the aza-nitrogen bridges and fused benzene rings to form MPc(-2) breaks this accidental degeneracy of the HOMO level. There is, therefore, reduced mixing between the Q and B excited states, and the Q band shifts to the red and gains significant intensity.<sup>36</sup> The expansion of the π-system through the addition of four additional fused

(58) Nonomura, T.; Kobayashi, N.; Tomura, T. *J. Porphyrins Phthalocyanines* **2000**, *4*, 538.

(59) Solovov, K. N.; Mashenkov, V. A.; Kachura, T. F. *Opt. Spectrosc.* **1969**, *27*, 24.



**Table 4.** Calculated Energies and Symmetries of the  $\pi$ -System MOs of a Model Complex Comprising the Inner Cyclic Polyene Ring of ZnPc (ZnCP), Zinc Tetraazaporphyrin (ZnTAP), ZnPc, Zn3B1N, Zn3BoN, Zncis2B2N, Zntrans2B2N, Zn1B3N, and ZnNpc(-2)<sup>a</sup>

$M_L$	ZnCP		ZnTAP		ZnPc(-2)		Zn3BoN		Zn3B1N		Zncis2B2N		Zntrans2B2N		Zn1B3N		ZnNpc(-2)	
	$D_{4h}$	E/eV	$D_{4h}$	E/eV	$D_{4h}$	E/eV	$C_s$	E/eV	$C_{2v}(III)$	E/eV	$C_{2v}(III)$	E/eV	$D_{2h}$	E/eV	$C_{2v}(III)$	E/eV	$D_{4h}$	E/eV
					3b <sub>2u</sub> *	3.844	21a''*	3.957	11a <sub>2</sub> *	3.957	12b <sub>2</sub> *	4.010	6a <sub>u</sub> *	4.016	13a <sub>2</sub> *	4.067	3b <sub>2u</sub> *	4.062
					5e <sub>g</sub> *	3.792	20a''*	3.892	10b <sub>2</sub> *	3.808	11b <sub>2</sub> *	3.897	6b <sub>3g</sub> *	4.013	12b <sub>2</sub> *	4.021	7e <sub>g</sub> *	4.036
						3.791	18a''*	3.748	10a <sub>2</sub> *	3.892	11a <sub>2</sub> *	3.929	6b <sub>2g</sub> *	3.822	11a <sub>2</sub> *	3.854		4.032
					2a <sub>1u</sub> *	3.678	19a''*	3.808	9a <sub>2</sub> *	3.748	10a <sub>2</sub> *	3.852	5a <sub>u</sub> *	3.788	12a <sub>2</sub> *	3.981	3a <sub>1u</sub> *	3.992
											8b <sub>2</sub> *	2.854	5b <sub>3g</sub> *	3.259	9a <sub>2</sub> *	3.119	6e <sub>g</sub> *	3.287
															11b <sub>2</sub> *	3.315	4b <sub>1u</sub> *	3.233
<b>8</b>	<b>1a<sub>2u</sub>*</b>	<b>2.260</b>	<b>1a<sub>2u</sub>*</b>	<b>2.310</b>	<b>2a<sub>2u</sub>*</b>	<b>2.961</b>	<b>16a''*</b>	<b>3.029</b>	<b>9b<sub>2</sub>*</b>	<b>3.054</b>	<b>10b<sub>2</sub>*</b>	<b>3.194</b>	<b>5b<sub>1u</sub>*</b>	<b>2.995</b>	<b>10b<sub>2</sub>*</b>	<b>2.996</b>	<b>4a<sub>2u</sub>*</b>	<b>3.024</b>
							17a''*	3.121	8a <sub>2</sub> *	3.121	8a <sub>2</sub> *	2.720	4a <sub>u</sub> *	3.083	8a <sub>2</sub> *	2.730	2a <sub>1u</sub> *	2.953
$\pm 7$	2e <sub>g</sub> *	2.385	3e <sub>g</sub> *	2.694	4e <sub>g</sub> *	2.945	15a''*	2.850	8b <sub>2</sub> *	2.903	9b <sub>2</sub> *	2.989	4b <sub>3g</sub> *	2.710	9b <sub>2</sub> *	2.756	5e <sub>g</sub> *	2.749
		2.385		2.693		2.946	14a''*	2.796	7a <sub>2</sub> *	2.800	9a <sub>2</sub> *	3.138	5b <sub>2g</sub> *	2.988	10a <sub>2</sub> *	3.145		2.747
			2b <sub>2u</sub> *	2.673	2b <sub>2u</sub> *	2.310	13a''*	2.302	6a <sub>2</sub> *	2.217	7b <sub>2</sub> *	2.168	3a <sub>u</sub> *	2.108	7a <sub>2</sub> *	2.109	2b <sub>2u</sub> *	2.261
			2b <sub>1u</sub> *		2b <sub>1u</sub> *	2.066	12a''*	2.086	7b <sub>2</sub> *	2.121	7a <sub>2</sub> *	2.168	4b <sub>1u</sub> *	2.176	8b <sub>2</sub> *	2.212	3b <sub>1u</sub> *	1.796
			2e <sub>g</sub> *	1.707	3e <sub>g</sub> *	1.677	11a''*	1.886	6b <sub>2</sub> *	1.743	6b <sub>2</sub> *	1.539	3b <sub>3g</sub> *	1.299	6b <sub>2</sub> *	1.469	4e <sub>g</sub> *	1.697
				1.704		1.676	10a''*	1.619	4a <sub>2</sub> *	1.205	6a <sub>2</sub> *	1.794	4b <sub>2g</sub> *	1.832	6a <sub>2</sub> *	1.665		1.693
			1a <sub>1u</sub> *	1.975	1a <sub>1u</sub> *	1.544	9a''*	1.497	5a <sub>2</sub> *	1.689	5a <sub>2</sub> *	1.392	2a <sub>u</sub> *	1.280	5a <sub>2</sub> *	1.437	1a <sub>1u</sub> *	1.124
					1a <sub>2u</sub> *	0.787	8a''*	1.101	5b <sub>2</sub> *	1.380	5b <sub>2</sub> *	1.331	3b <sub>1u</sub> *	1.488	5b <sub>2</sub> *	1.176	2a <sub>2u</sub> *	1.651
					2e <sub>g</sub> *	0.736	7a''*	0.749	4b <sub>2</sub> *	0.874	4b <sub>2</sub> *	0.907	2b <sub>3g</sub> *	1.457	4b <sub>2</sub> *	0.831	3e <sub>g</sub> *	1.189
						0.734	6a''*	0.667	3a <sub>2</sub> *	0.729	4a <sub>2</sub> *	1.050	3b <sub>2g</sub> *	0.728	4a <sub>2</sub> *	1.052		1.189
																	2e <sub>g</sub> *	0.023
											2a <sub>2</sub> *	-0.259	2b <sub>2g</sub> *	0.001	2a <sub>2</sub> *	0.001		0.019
							5a''*	0.573	2b <sub>2</sub> *	-0.208	2b <sub>2</sub> *	-0.130	2b <sub>1u</sub> *	-0.199	3b <sub>2</sub> *	-0.102	1a <sub>2u</sub> *	-0.024
$\pm 6$	1b <sub>2u</sub> *	1.644	1b <sub>2u</sub> *	1.180	1b <sub>2u</sub> *	0.098	4a''*	0.337	3b <sub>2</sub> *	0.507	3a <sub>2</sub> *	0.678	2b <sub>1u</sub> *	0.505	2b <sub>2</sub> *	-0.283	1b <sub>1u</sub> *	0.147
	1b <sub>1u</sub> *	0.763	1b <sub>1u</sub> *	-0.266	1b <sub>1u</sub> *	0.359	3a''*	-0.041	2a <sub>2</sub> *	0.107	3b <sub>2</sub> *	0.115	1a <sub>u</sub> *	0.125	3a <sub>2</sub> *	0.140	1b <sub>2u</sub> *	-0.295
$\pm 5$	1e <sub>g</sub> *	-0.989	1e <sub>g</sub> *	-1.843	1e <sub>g</sub> *	-1.495	2a''*	-1.447	1b <sub>2</sub> *	-1.415	1b <sub>2</sub> *	-1.423	1b <sub>3g</sub> *	-1.348	1b <sub>2</sub> *	-1.349	1e <sub>g</sub> *	-1.426
		-0.987		-1.842		-1.493	1a''*	-1.603	1a <sub>2</sub> *	-1.504	1a <sub>2</sub> *	-1.434	1b <sub>2g</sub> *	-1.563	1a <sub>2</sub> *	-1.479		-1.426
$\pm 4$	1a <sub>1u</sub>	-5.328	1a <sub>1u</sub>	-6.264	1a <sub>1u</sub>	-5.534	1a''	-5.600	1a <sub>2</sub>	-5.459	1a <sub>2</sub>	-5.389	1a <sub>u</sub>	-5.353	1a <sub>2</sub>	-5.292	1a <sub>1u</sub>	-5.183
	1a <sub>2u</sub>	-8.431	1a <sub>2u</sub>	-8.365	1a <sub>2u</sub>	-8.318	3a''	-8.413	1b <sub>2</sub>	-8.305	2b <sub>2</sub>	-8.298	1b <sub>1u</sub>	-8.289	2b <sub>2</sub>	-8.280	1a <sub>2u</sub>	-7.908
							2a''	-7.665	2a <sub>2</sub>	-8.134	1b <sub>2</sub>	-8.044	2a <sub>u</sub>	-8.153	2a <sub>2</sub>	-8.018	1b <sub>2u</sub>	-7.908
																	1e <sub>g</sub>	-7.914
											2a <sub>2</sub>	-8.202	1b <sub>3g</sub>	-7.937	1b <sub>2</sub>	-7.921		-7.915
															3a <sub>2</sub>	-8.268	2a <sub>1u</sub>	-8.417
			1b <sub>1u</sub>	-9.211	1b <sub>1u</sub>	-8.780	6a''	-8.911	2b <sub>2</sub>	-8.658	4a <sub>2</sub>	-8.719	2b <sub>1u</sub>	-8.652	3b <sub>2</sub>	-8.534	1b <sub>1u</sub>	-8.551
			1e <sub>g</sub>	-9.241	1e <sub>g</sub>	-8.811	4a''	-8.603	3b <sub>2</sub>	-8.735	3b <sub>2</sub>	-8.675	2b <sub>3g</sub>	-8.782	4b <sub>2</sub>	-8.704	2e <sub>g</sub>	-8.573
				-9.246		-8.813	5a''	-8.843	3a <sub>2</sub>	-8.802	3a <sub>2</sub>	-8.596	1b <sub>2g</sub>	-8.598	4a <sub>2</sub>	-8.594		-8.578
					2e <sub>g</sub>	-9.021	7a''	-9.037	4b <sub>2</sub>	-9.001	6b <sub>2</sub>	-10.167	3b <sub>3g</sub>	-9.989	6b <sub>2</sub>	-9.962	3e <sub>g</sub>	-9.933
						-9.021	8a''	-9.215	6a <sub>2</sub>	-10.263	6a <sub>2</sub>	-10.261	2b <sub>2g</sub>	-8.897	6a <sub>2</sub>	-10.183		-9.933
					1b <sub>2u</sub>	-9.161	11a''	-9.928	4a <sub>2</sub>	-9.000	4b <sub>2</sub>	-8.918	3a <sub>u</sub>	-9.421	5a <sub>2</sub>	-9.076	2b <sub>2u</sub>	-10.296
			2a <sub>2u</sub>	-9.765	2a <sub>2u</sub>	-9.297	9a''	-9.284	5b <sub>2</sub>	-9.233	5b <sub>2</sub>	-9.158	3b <sub>1u</sub>	-9.178	5b <sub>2</sub>	-9.110	2a <sub>2u</sub>	-9.053
			2a <sub>1u</sub>		2a <sub>1u</sub>	-9.711	10a''	5a <sub>2</sub>	-9.526	5a <sub>2</sub>	-9.274	4a <sub>u</sub>	-10.478	7a <sub>2</sub>	-10.440	3a <sub>1u</sub>	-10.563	
$\pm 3$	1e <sub>g</sub>	-10.390	2e <sub>g</sub>	-10.399	3e <sub>g</sub>	-10.695	12a''	-10.600	6b <sub>2</sub>	-10.619	7b <sub>2</sub>	-10.793	4b <sub>3g</sub>	-10.932	7b <sub>2</sub>	-10.858	4e <sub>g</sub>	-10.779
	1e <sub>g</sub>	-10.392		-10.401		-10.698	13a''	-10.758	7a <sub>2</sub>	-10.893	7a <sub>2</sub>	-10.872	3b <sub>2g</sub>	-10.538	8a <sub>2</sub>	-10.811		-10.783
															8b <sub>2</sub>	-11.283	2b <sub>1u</sub>	-11.138
							14a''	-11.583	7b <sub>2</sub>	-11.614	8b <sub>2</sub>	-11.825	4b <sub>1u</sub>	-11.369	9b <sub>2</sub>	-11.819	3a <sub>2u</sub>	-11.678
																	5e <sub>g</sub>	-11.712
											8a <sub>2</sub>	-11.470	5b <sub>2g</sub>	-11.713	9a <sub>2</sub>	-11.710		-11.715
$\pm 2$	1b <sub>2u</sub>	-12.730	1b <sub>2u</sub>	-12.782	2b <sub>2u</sub>	-12.275	16a''	-12.875	8a <sub>2</sub>	-12.854	9b <sub>2</sub>	-12.845	5a <sub>u</sub>	-12.869	10a <sub>2</sub>	-12.883	3b <sub>2u</sub>	-12.852
	1b <sub>1u</sub>	-13.063	2b <sub>1u</sub>	-13.670	2b <sub>1u</sub>	-12.870	15a''	-12.517	8b <sub>2</sub>	-12.469	9a <sub>2</sub>	-12.565	5b <sub>1u</sub>	-12.720	10b <sub>2</sub>	-12.693	3b <sub>1u</sub>	-13.709
					4e <sub>g</sub>	-13.088	17a''	-13.124	9b <sub>2</sub>	-13.087	10b <sub>2</sub>	-13.018	6b <sub>3g</sub>	-14.432	11b <sub>2</sub>	-13.964	6e <sub>g</sub>	-14.410
						-13.089	18a''	-13.336	9a <sub>2</sub>	-13.115	10a <sub>2</sub>	-14.163	6b <sub>2g</sub>	-13.602	11a <sub>2</sub>	-14.427		-14.411
					3a <sub>2u</sub>	-13.292	19a''	-14.070	10b <sub>2</sub>	-14.420	11b <sub>2</sub>	-14.665	6b <sub>1u</sub>	-14.240	12b <sub>2</sub>	-14.686	4a <sub>2u</sub>	-14.847
					3b <sub>1u</sub>	-15.157	20a''	-15.289	11b <sub>2</sub>	-15.270	11a <sub>2</sub>	-15.327	7b <sub>1u</sub>	-15.401	13b <sub>2</sub>	-15.419	4b <sub>1u</sub>	-15.532
$\pm 1$	2e <sub>g</sub>	-14.898	3e <sub>g</sub>	-15.040	5e <sub>g</sub>	-16.085	21a''	-16.110	10a <sub>2</sub>	-16.070	12a <sub>2</sub>	-16.091	7b <sub>2g</sub>	-16.21	12a <sub>2</sub>	-16.197	7e <sub>g</sub>	-16.173
		-14.898		-15.046		-16.086	22a''	-16.153	12b <sub>2</sub>	-16.149	12b <sub>2</sub>	-16.176	6b <sub>3g</sub>	-16.052	14b <sub>2</sub>	-16.174		-16.175
<b>0</b>	2a <sub>2u</sub>	-16.256	3a <sub>2u</sub>	-16.151	4a <sub>2u</sub>	-17.101	23a''	-17.132	13b <sub>2</sub>	-17.111	13a <sub>2</sub>	-17.118	8b <sub>1u</sub>	-17.104	15b <sub>2</sub>	-17.124	5a <sub>2u</sub>	-17.102

<sup>a</sup> The rows are ordered according to the nodal and antinodal patterns that are observed for each complex relative to the high symmetry ZnPc(-2) parent complex. The orbital angular momentum associated with the orbitals of the inner cyclic polyene model complex are shown in the left-hand column. The orbitals with the corresponding nodal patterns are highlighted in bold. Italics are used to indicate degenerate orbitals under  $D_{4h}$  symmetry.

benzenes to form Nc(-2) results in a further increase in Q band intensity and a further red shift of the Q band center.<sup>60</sup> As would be expected, the  $A_1/D_0$  ratio obtained from MCD spectroscopy decreases on going from ZnPc(-2) to ZnNc(-2) as the separation between the 1a<sub>1u</sub> and 1a<sub>2u</sub> levels increases. The energy differences of the Q band center of ZnPc(-2) and Zncis2B2N (934 cm<sup>-1</sup>) and that of Zncis2B2N and ZnNc(-2) (876 cm<sup>-1</sup>) are reasonably similar, Figure 5.

(60) Kobayashi, N.; Nakijima, S.; Osa, T. *Inorg. Chim. Acta* **1993**, *210*, 131.

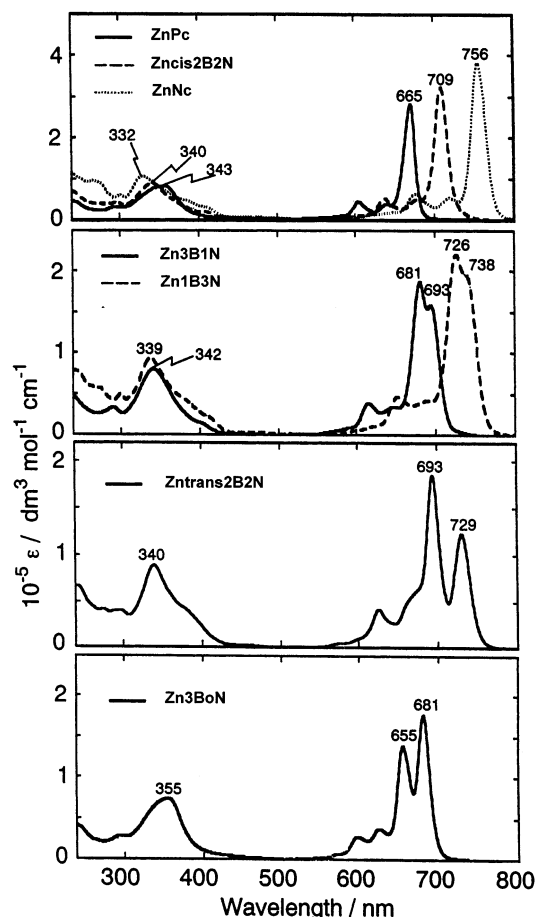
There is a slight red shift of the Q bands of Zn3B1N at 682 and 697 nm, relative to that of ZnPc(-2) at 670 nm,<sup>3a</sup> Table 4. In contrast, there is a slight blue shift of the Q band components of Zn1B3N at 725 and 743 nm (using the MCD B term band centers) relative to the Q band of ZnNc(-2) at 750 nm.<sup>3a</sup>

The ZINDO calculations successfully predict the magnitude of the energy gap between the LUMO and HOMO and the associated red shift of the Q band as naphthalene groups are added to the Pc ring, Figure 4. The predicted splitting of

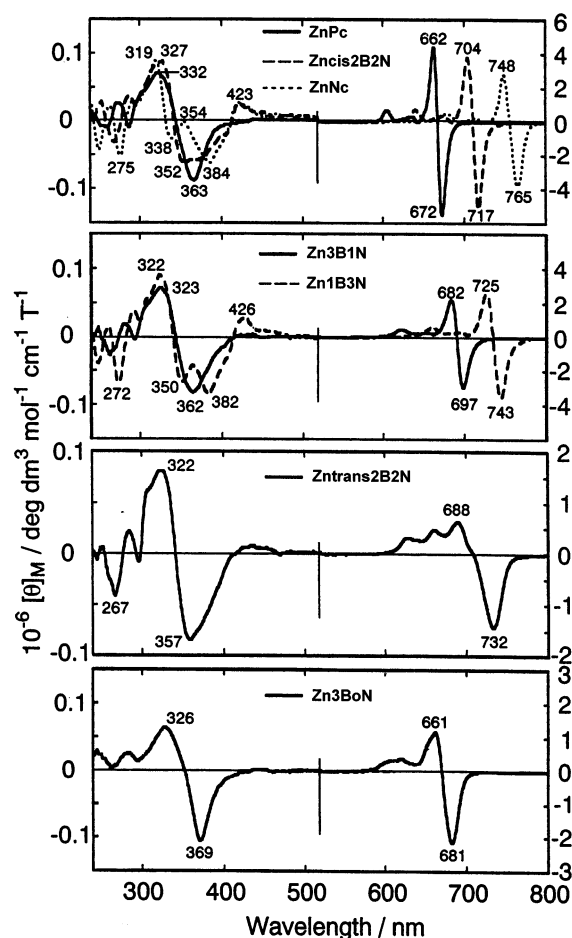
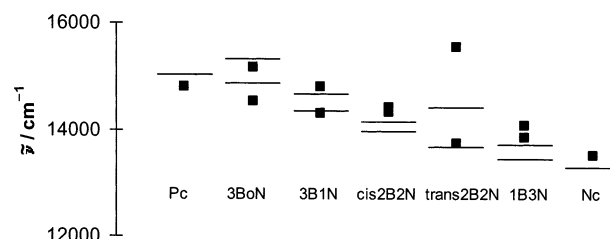
**Table 5.** Calculated Energies and Symmetries of the Primarily Aza-Bridging and Pyrrole Nitrogen Lone Pair Based  $\sigma$ -MOs in the  $-11$  to  $3$  eV Range for a Model Complex Comprising the Inner Cyclic Polyene Ring of ZnPc (ZnCP), Zinc Tetraazaporphyrin (ZnTAP), ZnPc, Zn3B1N, Zn3BoN, Zncis2B2N, Zntrans2B2N, Zn1B3N, and ZnNpc(-2)<sup>a</sup>

ZnPc(-2) $D_{4h}$		Zn3BoN $C_s$		Zn3B1N $C_{2v}(\text{III})$		Zncis2B2N $C_{2v}(\text{III})$		Zntrans2B2N $D_{2h}$		Zn1B3N $C_{2v}(\text{III})$		ZnNpc(-2) $D_{4h}$	
$1b_{1g}$	-9.532	$1a'$	-9.445	$1b_1$	-9.493	$1a_1$	-9.447	$1b_{1g}$	-9.493	$1b_1$	-9.411	$1b_{1g}$	-9.488
$1b_{2g}$	-9.766	$2a'$	-9.750	$1a_1$	-9.746	$1b_1$	-9.728	$1a_u$	-9.746	$1a_1$	-9.739	$1b_{2g}$	-9.759
$1e_u$	-10.013	$3a'$	-9.945	$2a_1$	-9.983	$2b_1$	-9.988	$2b_{3u}$	-9.983	$2a_1$	-9.948	$1e_u$	-9.958
$1e_u$	-10.018	$4a'$	-10.025	$2b_1$	-10.035	$2b_1$	-10.017	$2b_{2u}$	-10.035	$2b_1$	-10.001	$1e_u$	-9.958
$1a_{1g}$	-10.817	$5a'$	-10.713	$3a_1$	-10.831	$3a_1$	-10.863	$1a_g$	-10.493	$3a_1$	-10.780	$1a_{1g}$	-10.740

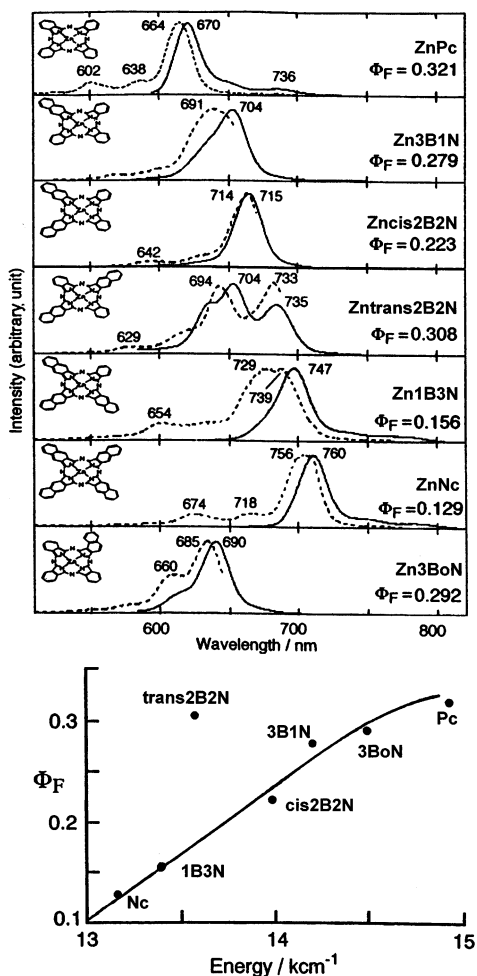
<sup>a</sup> The rows are ordered according to the nodal and antinodal patterns that are observed for each complex relative to the high symmetry ZnPc(-2) parent complex. Italics are used to indicate degenerate orbitals under  $D_{4h}$  symmetry.

**Figure 5.** Electronic absorption spectra of ZnPc(-2) 1, Zn3B1N 2, Zncis2B2N 3, Zntrans2B2N 4, Zn1B3N 5, ZnNc(-2) 6, and Zn3BoN 7 in THF.

the two Q band components also closely matches the experimental values, Figure 7. The splitting observed in the Q bands of Zntrans2B2N at 688 and 732 nm is significantly greater than is the case with Zncis2B2N where bands are observed at 704 and 717 nm despite the fact that the Zntrans2B2N complex is the higher symmetry complex. The perturbation to the structure of ZnPc(-2) results in different ring substitution patterns along the  $x$ - and  $y$ -axes in the case of the  $D_{2h}$  Zntrans2B2N isomer while the structure is identical along the  $x$ - and  $y$ -axes in the case of Zncis2B2N, Figure 2. The spectrum in the Q region of Zncis2B2N is therefore very similar to that of ZnPc(-2) and ZnNc(-2) as the Q bands of the  $C_{2v}(\text{III})$  Zncis2B2N isomer form a pseudo-A term because the excited states associated with the  $1a_2 \rightarrow 1a_2^*$  and  $1a_2 \rightarrow 1b_2^*$  one-electron transitions are

**Figure 6.** MCD spectra of ZnPc(-2) 1, Zn3B1N 2, Zncis2B2N 3, Zntrans2B2N 4, Zn1B3N 5, ZnNc(-2) 6, and Zn3BoN 7 in THF.**Figure 7.** Predicted energies of the  $Q_{00}$  bands (-) from ZINDO/s plotted against the experimentally observed values (■) obtained from spectral deconvolution, Tables S4-S8.

nearly degenerate. The spectra of Zn3B1N and Zn3BoN are at first glance very similar to those of ZnPc(-2),<sup>3a,e</sup> but the reduced symmetry results in a slight splitting of the Q band because of the lifting of the orbital degeneracy of the  $1e_g^*$



**Figure 8.** Fluorescence emission and excitation spectra of ZnPc(−2) 1, Zn3B1N 2, Zncis2B2N 3, Zntrans2B2N 4, Zn1B3N 5, ZnNc(−2) 6, and Zn3BoN 7 in deaerated CHCl<sub>3</sub>.

LUMO. There is a slightly greater splitting between the Q components of Zn3BoN because of the impact of the reduced molecular symmetry, Table 4, but Q band positions of Zn3BoN are closer to that of ZnPc(−2) than is the case with Zn3B1N. These results are consistent with the existing spectral data, which have suggested that oblique substitution or fusion has a relatively minor impact on the wavelength of the Q band.<sup>19a,61</sup> The MCD spectrum of Zn3BoN appears to be almost identical to that of Zn3B1N in the B region while that of Zn1B3N is considerably more complex with significantly greater intensity within the spectral envelope in the 300–450 nm region. Many of the new MOs which result from  $\pi$ -system expansion have energies close to those of the four frontier orbitals of the inner cyclic polyene, Table 4 and Figure 4, so additional one-electron transitions can be expected to contribute to band intensity in this region.

Figure 8 displays the S1 fluorescence emission and excitation spectra of the complexes studied, together with their  $\Phi_F$  values. The observed Stokes shifts were relatively small in all cases, and Zntrans2B2N is the only complex to show two well-resolved Q bands in the emission and

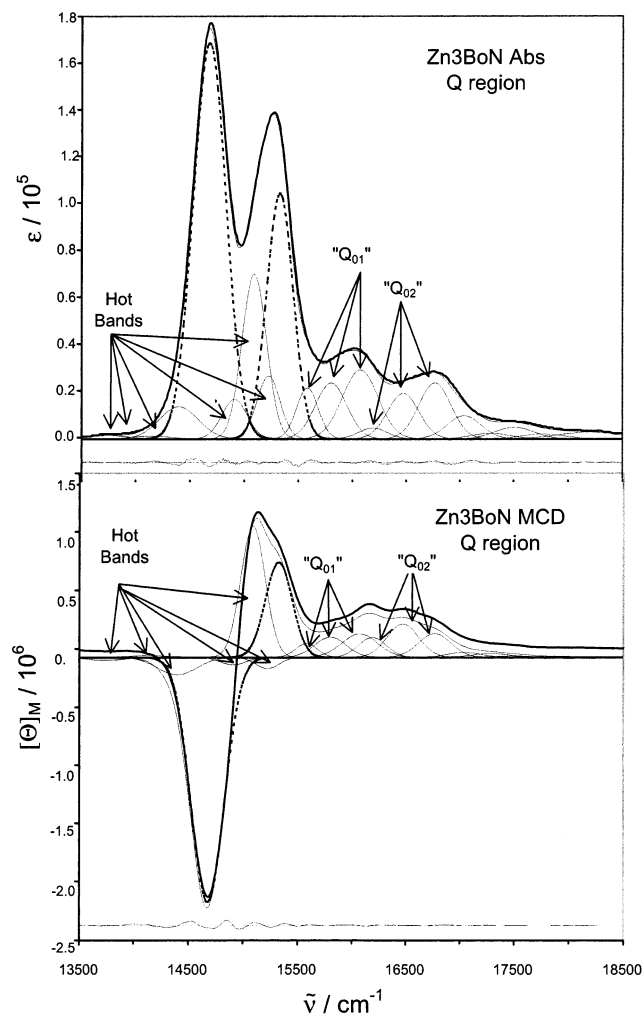
excitation spectra. A deviation from mirror symmetry similar to that reported previously for MPc(−2)<sup>36,62</sup> was observed in the emission and excitation spectra of all the complexes and was particularly marked in the case of the  $C_{2v}$ (III) symmetry complexes. When the  $\Phi_F$  value is plotted against the wavelength of the Q band, Figure 8, a clear linear relationship is observed. This strongly suggests that non-radiative decay becomes easier as the energy gap decreases between the HOMO and the LUMO. The  $D_{2h}$  Zntrans2B2N complex where there is the largest energy separation in the  $x$ - and  $y$ -polarized Q00 bands represents the only significant outlier from this trend, which indicates that this fused ring substitution pattern has the most profound impact on the electronic properties of the  $\pi$ -system. Emission spectra were recorded at several excitation wavelengths.

**(b) Band Deconvolution: The Q and  $n \rightarrow \pi^*$  Spectral Regions.** The band deconvolution analyses of the Q region of the absorption and MCD spectra of Zn3BoN, Zn3B1N, Zncis2B2N, Zntrans2B2N, and Zn1B3N provide direct quantitative spectroscopic data on the effect of the asymmetry on the electronic structure, Figures 9 and 10 and Figures S1–S3 (Supporting Information). The two major electronic bands arising from the split Q00 transition are shown in bold. The band deconvolution analysis of spectra obtained on a vitrified ZnPc(−2) solution at cryogenic temperatures, Figure 11, led to the identification of a repeating pattern within the envelope of vibrational bands associated with the Q transition.<sup>36</sup> An  $a_{1g}$  vibration combines with the first set of vibrations so that repeating “Q01” and “Q02” envelopes of vibrational bands dominate the spectrum between the Q00 band and the 607 nm band. The energies obtained from our deconvolution analysis for the major vibrational bands were in close agreement with those identified by Huang et al. from fluorescence emission and excitation spectra in low temperature argon matrices.<sup>62</sup> In the case of ZnPc(−2), a mismatch in the wavelengths of the absorption and MCD bands at ca. 600 nm and the absence of mirror symmetry in the fluorescence and excitation spectra which is not observed in the case of MP(−2) complexes were found to be consistent with a weak  $n \rightarrow \pi^*$  band associated with an MO located primarily on the aza-nitrogen lone pair electrons being the origin for a second weaker set of vibrational bands within the Q spectral envelope.<sup>36,62</sup> The components identified by this deconvolution analysis should be present in the Q spectral region of all phthalocyanine complexes.

A complex envelope of vibrational bands, similar to that observed for MPc(−2), Figure 11, is observed to the blue of the Q band in each of the 5 spectral data sets, Figures 9 and 10 and Figures S1–S3 (Supporting Information). Some minor hot bands associated with the Boltzmann population distribution across the vibrational levels associated with the ground states are observed to the red of the Q band like those observed previously in the case of ZnPc(−2), Figure 12. These hot bands are eliminated in data recorded on vitreous solutions at cryogenic temperatures, Figure 11.<sup>36</sup> Band

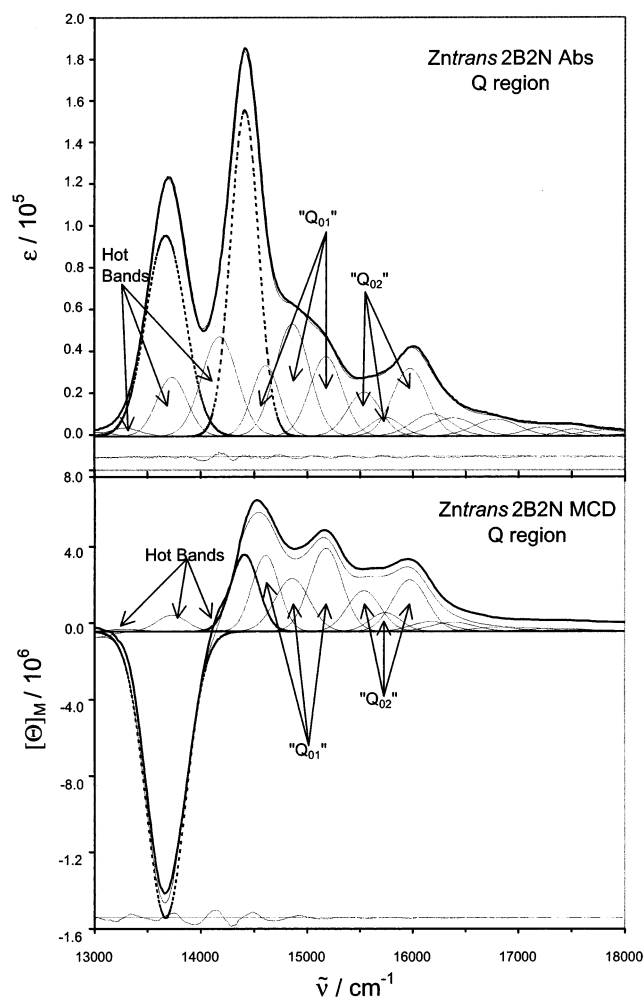
(61) Luk'yanyets, E. A. *Electronic Spectra of Phthalocyanines and Related Compounds*; NIOPIK: Moscow, 1989.

(62) (a) Huang, T. H.; Rieckhoff, K. E.; Voight, E. M. *J. Chem. Phys.* **1982**, *77*, 3424. (b) Huang, T. H.; Rieckhoff, K. E.; Voight, E. M. *J. Phys. Chem.* **1981**, *85*, 3322.



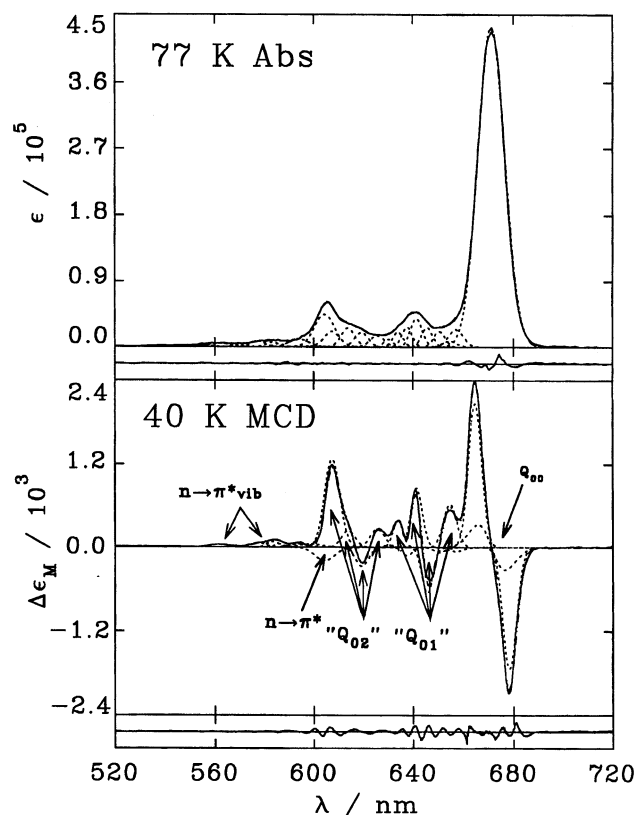
**Figure 9.** Spectral band deconvolution analysis of the Q spectral region of Zn3BoN. The band fitting parameters can be found in Table S4. Major electronic bands assigned in Table 2 to the Q transition, Figure 3, are shown in bold. The line showing the experimental MCD spectrum has been shifted positive slightly for visual purposes only.

broadening due to solvation effects and the splitting of the Q band as a result of reduced symmetry further complicate the spectra and result in a very high degree of band overlap. The rigid shift assumption under which the bandwidths in both the UV–vis and MCD spectra can be assumed to be identical is problematic when hot bands are present.<sup>30</sup> A repeating pattern of “Q01” and “Q02” envelopes can, however, still be seen in the SIMPFIT spectral fits of all five complexes consisting of weak positive overlapping B terms as was observed previously in the case of ZnPc(–2), Figure 11. The analysis provides important new information. The B terms associated with some of the vibrational bands of Zn3BoN and Zntrans2B2N, the complexes where there is the largest splitting of the main Q bands, are more intense than the B term associated with the higher energy component of the symmetry-split Q absorption band, Figures 9 and 10. The higher energy of the symmetry-split Q bands therefore lies to the blue of some of the vibrational bands associated with the lower energy band resulting in a spectrum markedly different from that of ZnPc(–2) because B term intensity is related to the field-induced mixing of close-lying states.



**Figure 10.** Spectral band deconvolution analysis of the Q spectral region of Zntrans2B2N. The band fitting parameters can be found in Table S7. Major electronic bands assigned in Table 3 to the Q transition, Figure 3, are shown in bold. The line showing the experimental MCD spectrum has been shifted positive slightly for visual purposes only.

On the blue side of the Q spectral envelope, the calculations identify a set of bands with widths that do not fit the pattern of positive and negative B terms in the “Q01” and “Q02” vibronic progression, Figure 11. These bands are assigned to the weak  $n \rightarrow \pi^*$  transition observed in the case of ZnPc(–2).<sup>36</sup> There is a significant deviation from mirror symmetry similar to that originally observed by Huang in the fluorescence emission and excitation spectra of ZnPc(–2),<sup>36,62</sup> Figure 8. If this assignment is valid, the ZINDO calculations must significantly underestimate the energy of the  $\sigma$ -MOs based primarily on the aza-nitrogen lone pair orbitals.<sup>54</sup> An MO calculation by Schaffer et al.<sup>28</sup> predicted that the aza-nitrogen lone pair MOs lie in the same energy range as the  $a_{1u}$  HOMO orbital while the same MOs lie about 3 eV below the  $a_{1u}$  HOMO in a calculation by Henriksson et al.<sup>29a</sup> Unlike the MO calculations, the experimental data obtained from optical spectroscopy are unambiguous and fixed, and the assignment of these bands to a weak  $n \rightarrow \pi^*$  transition appears to provide the only explanation that can fully account for all the features observed within the spectrum. An anomalous, intense B term in the spectrum of [MPc(–1)]<sup>+</sup> at ca. 500 nm to the blue of the Q band in a

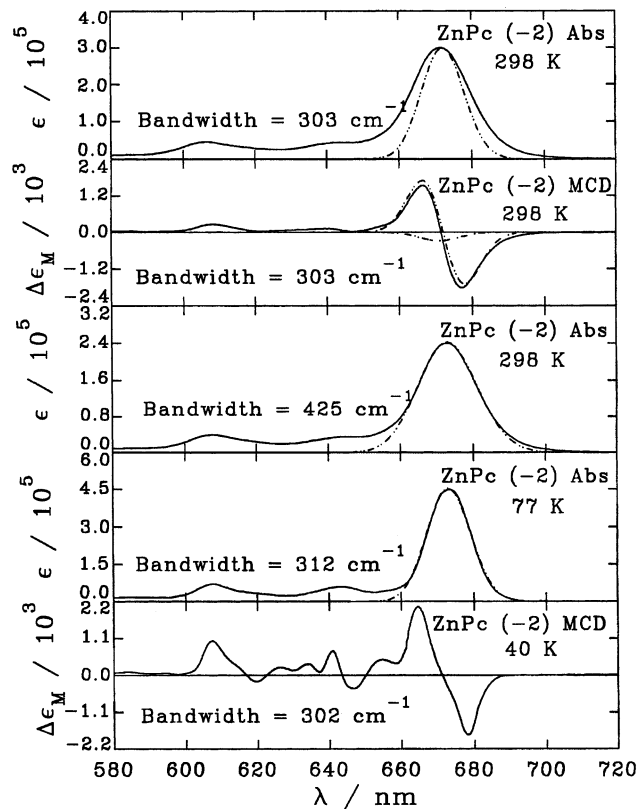


**Figure 11.** Spectral deconvolution analysis of the Q band region of the absorption and MCD spectra of  $(\text{CN}^-)\text{ZnPc}(-2)$ .<sup>36</sup> Identical bandwidths and centers are used to fit the MCD and absorption spectra recorded at 40 and 77 K on a 5:2 dimethyl formamide/dimethyl acetamide vitrified solution. The solid, dashed and dotted lines represent the experimental and calculated spectral data and the component deconvoluted B terms, respectively. All the bands used in this analysis with the exception of the main  $Q_{00}$  band are B terms. The  $n \rightarrow \pi^*$  bands near 605 nm and its associated vibronic progression are indicated. The residual of the fit plotted at the bottom of the plot shows the difference between the experimental and calculated MCD data. Reproduced with permission from ref 36. Copyright 1995 American Chemical Society.

spectrum dominated by A terms provides further evidence to support this assignment.<sup>36,54</sup>

### (c) Band Deconvolution: The B1/B2 Spectral Region.

Spectral deconvolution of absorption and MCD spectra has indicated previously that there are two separate A terms under Gouterman's B band in the spectrum of  $\text{MPc}(-2)$  complexes.<sup>3a,e,31–33,36</sup> ZINDO calculations indicate that this is due to mixing of the B transition with a one-electron transition involving the  $1b_{1u}$  orbital.<sup>54</sup> A very weak A term sometimes observed slightly to the red of the B1 band has been assigned to a second  $\pi \rightarrow \pi^*$  transition between the  $a_{1u}$  HOMO and a degenerate orbital located primarily on the fused benzene periphery of the ring system that is not present in the case of  $\text{MP}(-2)$ .<sup>54</sup> The band deconvolution analyses of the B region of the absorption and MCD spectra of  $\text{Zn3BoN}$ ,  $\text{Zn3B1N}$ ,  $\text{Zncis2B2N}$ ,  $\text{Zntrans2B2N}$ , and  $\text{Zn1B3N}$  are shown in Figures 13 and 14. The B1, B2, and to a lesser extent the second  $\pi \rightarrow \pi^*$  one-electron transitions are predicted to contribute intensity to the most prominent bands observed between 26000 and 32000  $\text{cm}^{-1}$  for all five complexes, Tables 1–3 and Tables S1–S3 (Supporting Information). The B1 bands remain the most intense UV–vis absorption bands because of the high degree of orbital overlap between

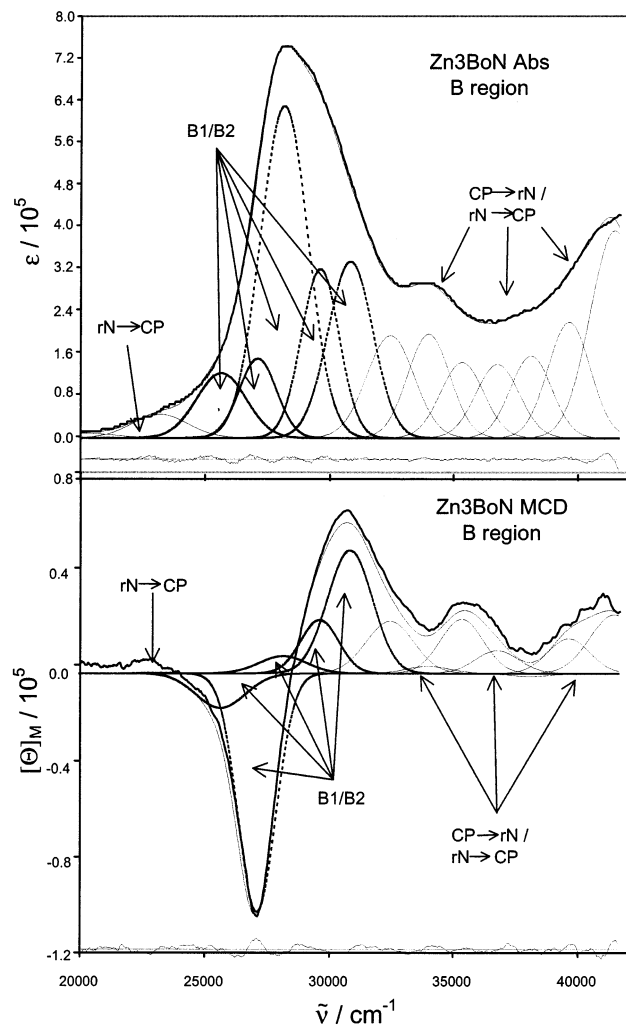


**Figure 12.** Temperature dependence in the Q band region of the UV–vis absorption and MCD spectra of  $\text{ZnPc}(-2)$  in liquid and vitrified solutions. The Gaussian curves illustrate the bandwidth discrepancy in the room temperature data. Reproduced with permission from ref 36. Copyright 1995 American Chemical Society.

the frontier orbitals associated with the cyclic polyene. This allowed  $\Delta M_L = \pm 1$  transition of the inner cyclic polyene also dominates the MCD spectra because of the OAM associated with these orbitals, Figures 13 and 14 and S4–S6 (Supporting Information). Comparison with the calculated absorption spectra, Tables 1–3 and S1–S3 (Supporting Information), indicates that the B1 band tends to be ca. 5000  $\text{cm}^{-1}$  to the red of the values predicted in the ZINDO calculations. The second  $\pi \rightarrow \pi^*$  band shifts to the red as the orbitals corresponding to the  $2e_g^*$  orbital of  $\text{ZnPc}(-2)$ , which are associated primarily with the four peripheral benzo-groups, lie at higher energy relative to the HOMO orbital in most of the fused-ring substituted complexes, Table 4. The weaker bands to the blue of the B1 bands, Figures 13 and 14 and S4–S6 (Supporting Information), arise primarily from one-electron transitions involving both the inner cyclic polyene ring and the new peripheral fused rings. The original N/L/C band assignment terminology, which has previously been applied to the experimental spectra of  $\text{MPc}(-2)$ ,<sup>54</sup> is, therefore, clearly no longer directly applicable when there is an asymmetric extension to the  $\pi$ -system through the addition of peripheral fused rings.

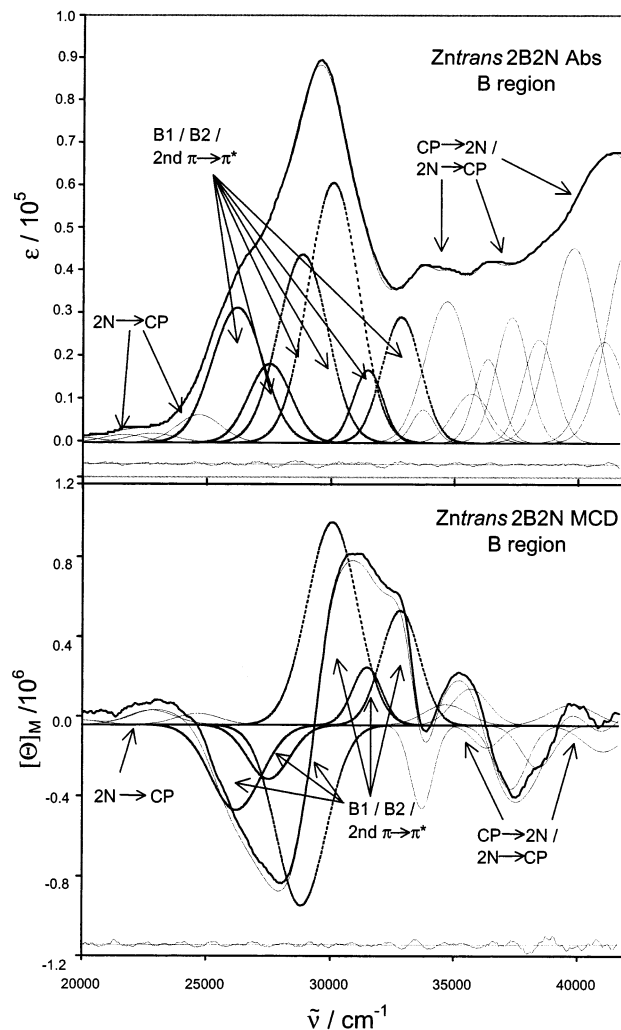
## Conclusions

Although the assignment of all the NMR signals for the asymmetric fused-ring-substituted complexes was relatively straightforward, some of the shifts observed for  $\text{Zntrans2B2N}$  and  $\text{Zn3BoN}$  do not correspond with the standard ring current



**Figure 13.** Spectral band deconvolution analysis of the B1/B2 spectral region of Zn3BoN. The band fitting parameters can be found in Table S9. Major electronic bands assigned in Table 2 to the B1 and B2 transitions, Figure 3, are shown in bold. The line showing the experimental MCD spectrum has been shifted positive slightly for visual purposes only.

interpretation developed previously to account for the NMR spectra of MP(-2) and MPc(-2) complexes. The optical spectra of Zntrans2B2N and Zn3BoN are also markedly different from those of Zn1B3N, Zncis2B2N, and Zn3B1N because there is a more marked symmetry-induced splitting in the Q spectral region. The bands associated primarily with the Q and B1 one-electron transitions of the 16-atom 18- $\pi$ -electron inner cyclic polyene remain the dominant features in the spectra of all five complexes. The direct relevance of band assignment models developed previously for MPc(-2) complexes is, however, more limited in the UV region than was the case with partially aza-substituted complexes.<sup>37</sup> ZINDO/s calculations indicate the presence of additional MOs associated primarily with the additional peripheral fused benzene rings at energies similar to the  $1a_{2u}$  orbital of MPc(-2). The spectral impact of these orbitals is most pronounced in the spectrum of Zn1B3N. The  $\Phi_F$  values, NMR spectroscopy, and spectral band deconvolution analyses of



**Figure 14.** Spectral band deconvolution analysis of the B1/B2 spectral region of Zntrans2B2N. The band fitting parameters can be found in Table S12. Major electronic bands assigned in Table 3 to the B1 and B2 transitions, Figure 3, are shown in bold. The line showing the experimental MCD spectrum has been shifted positive slightly for visual purposes only.

the optical spectra identify Zntrans2B2N as the complex with the most substantially modified electronic properties relative to ZnPc(-2). Complexes based on this asymmetric fused ring pattern therefore probably have the most potential for use in new industrial applications.

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**Supporting Information Available:** Calculated UV-vis spectra (Tables S1–S3), spectral band deconvolution analyses of the Q (Figures S1–S3) and B1/B2 (Figures S4–S6) spectral regions of Zn3B1N, Zncis2B2N, and Zn1B3N. Band fitting parameters from SIMPFIT deconvolution analyses of the Q and B1/B2 spectral regions (Tables S4–S13) of Zn3BoN, Zn3B1N, Zncis2B2N, Zntrans2B2N, and Zn1B3N. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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